

# Sudipfest

UCLA

April 4-6, 2019

*Intertwined Charge, Nematic,  
Magnetic and Superconducting Order  
in Fe-Based Superconductors*

Robert J. Birgeneau

UC Berkeley

1. **DYNAMICS OF THE DISSIPATIVE 2-STATE SYSTEM**

By: LEGGETT, AJ; CHAKRAVARTY, S; DORSEY, AT; et al.  
REVIEWS OF MODERN PHYSICS Volume: 59 Issue: 1 Pages: 1-85 Published: JAN 1987

2. **TWO-DIMENSIONAL QUANTUM HEISENBERG-ANTIFERROMAGNET AT LOW-TEMPERATURES**

By: CHAKRAVARTY, S; HALPERIN, BI; NELSON, DR  
PHYSICAL REVIEW B Volume: 39 Issue: 4 Pages: 2344-2371 Published: FEB 1 1989

3. **Hidden order in the cuprates**

By: Chakravarty, S; Laughlin, RB; Morr, DK; et al.  
PHYSICAL REVIEW B Volume: 63 Issue: 9 Article Number: 094503 Published: MAR 1 2001

4. **WEAK LOCALIZATION - THE QUASI-CLASSICAL THEORY OF ELECTRONS IN A RANDOM POTENTIAL**

By: CHAKRAVARTY, S; SCHMID, A  
PHYSICS REPORTS-REVIEW SECTION OF PHYSICS LETTERS Volume: 140 Issue: 4 Pages: 193-236 Published: JUL 1986

5. **QUANTUM FLUCTUATIONS IN THE TUNNELING BETWEEN SUPERCONDUCTORS**

By: CHAKRAVARTY, S  
PHYSICAL REVIEW LETTERS Volume: 49 Issue: 9 Pages: 681-684 Published: 1982

135	133	146	137	17	3540
18	21	32	20	4	1207
32	22	16	11	2	841
1	8	6	7	1	357
3	7	1	4	2	335

## Group Members

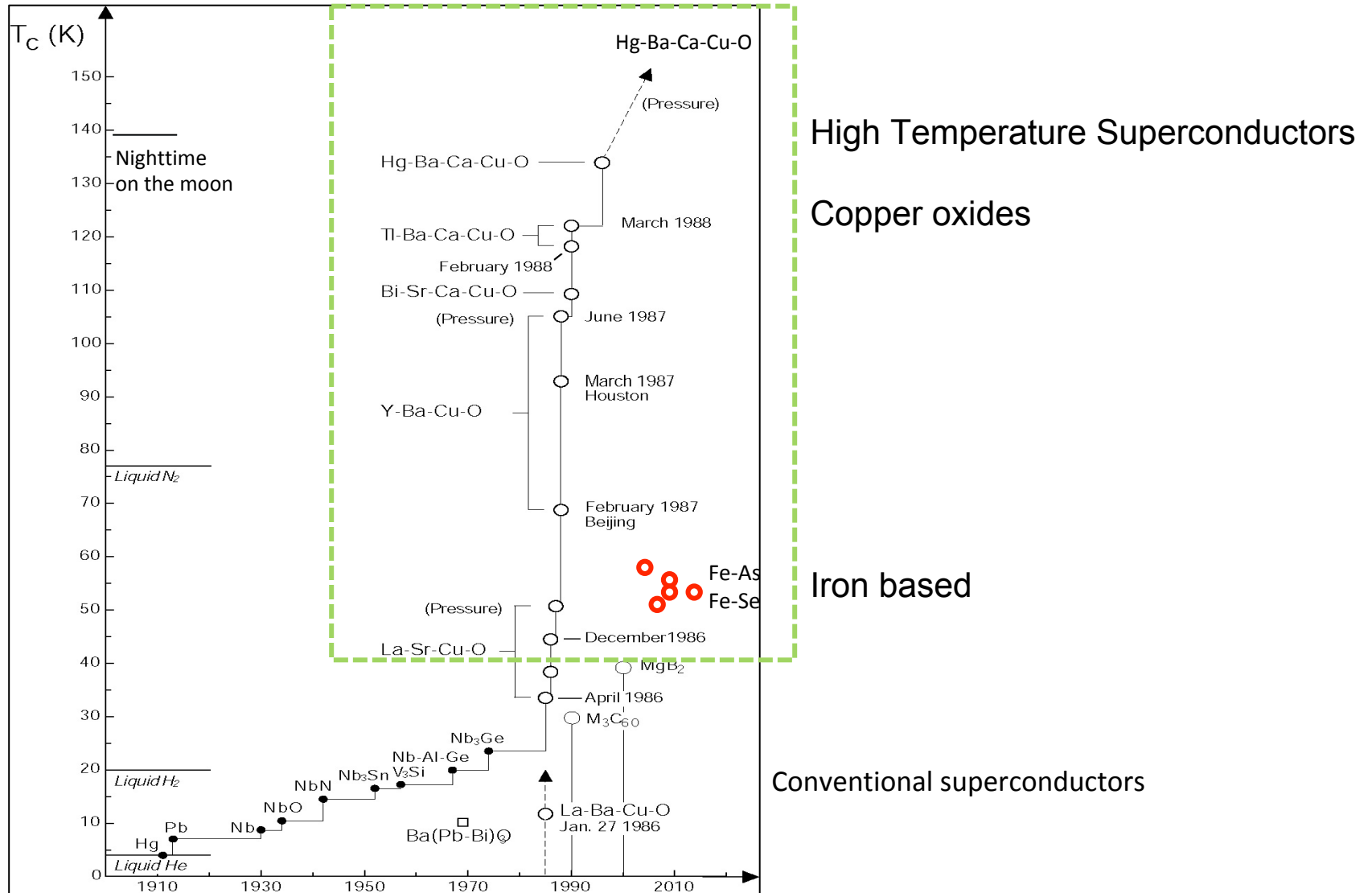
Xiang Chen  
Ben Frandsen\*  
Alex Frano\*  
Yu He  
Min Gyu Kim\*  
Yu Song  
Meng Wang\*  
Shan Wu  
Zhijun Xu\*  
Ming Yi\*  
Jun Zhao\*

\*Recent past members

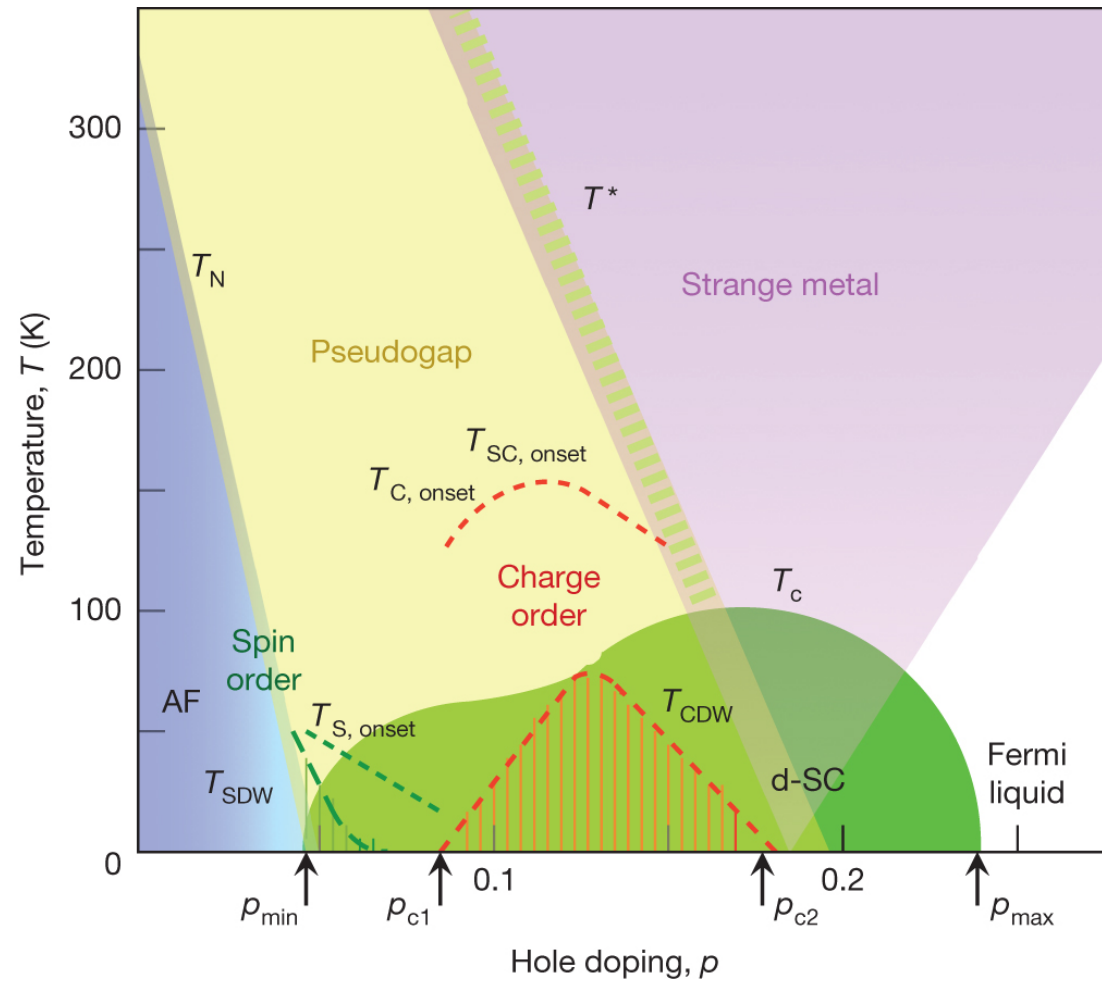
## Collaborators

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Byron Freelon  
Bernhard Keimer  
Alessandra Lanzara  
Dung Hai Lee  
Z.-X. Shen  
John Tranquada  
Stephen Wilson

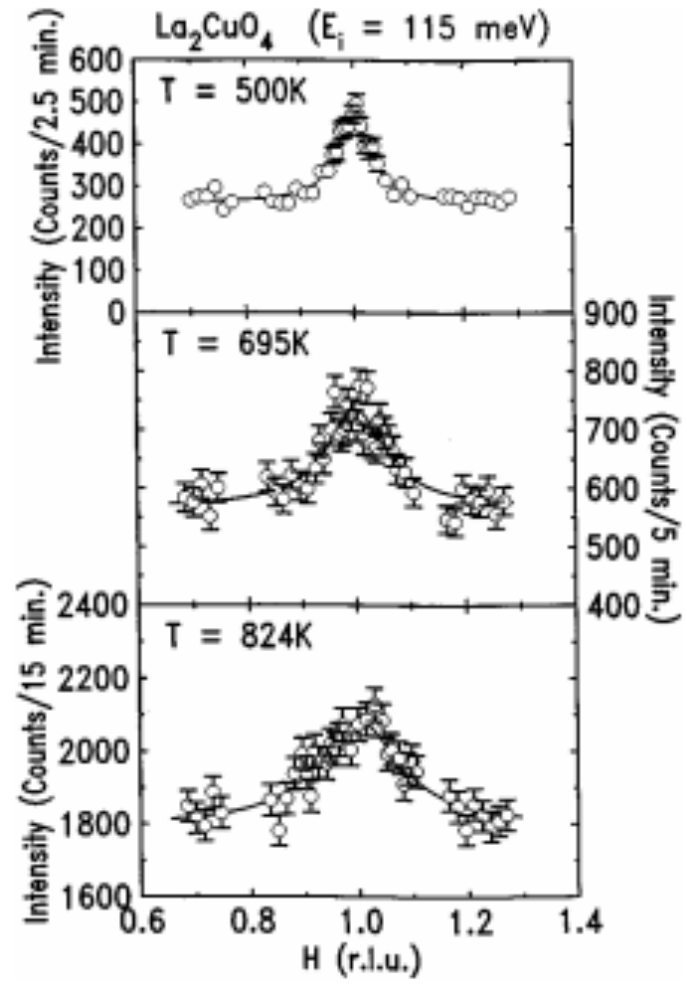
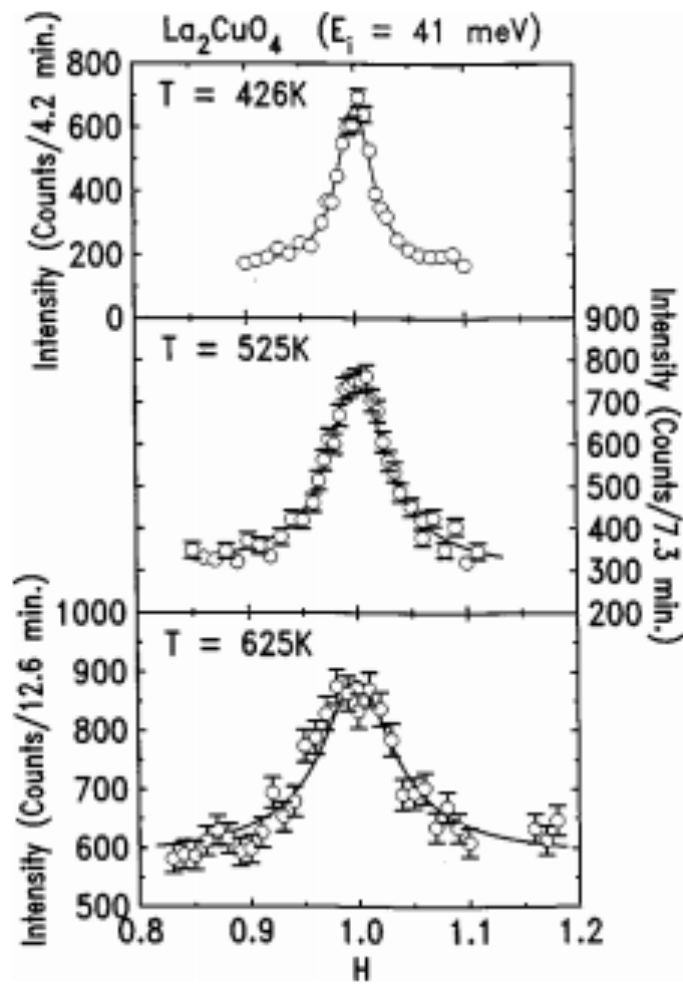
# History of Superconductivity



# Phase diagram.



B Keimer *et al.* *Nature* **518**, 179-186 (2015) doi:10.1038/nature14165



# La<sub>2</sub>CuO<sub>4</sub>: A Model Two Dimensional S = 1/2 Heisenberg Antiferromagnet

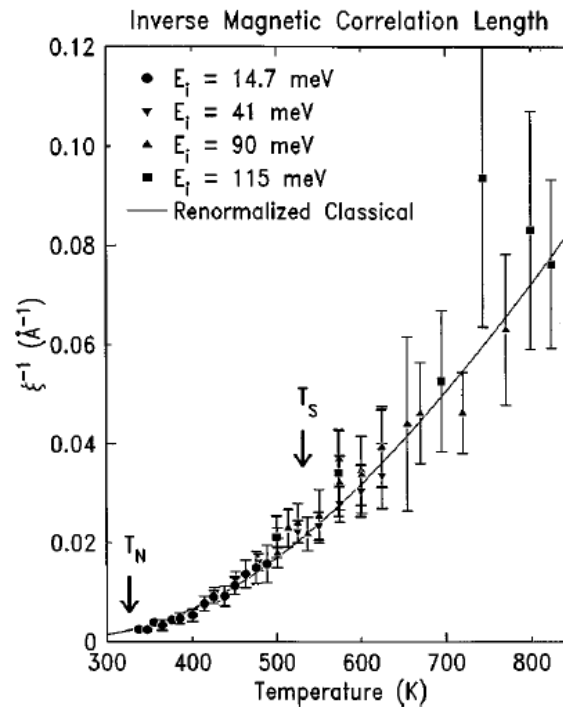


FIG. 5. Inverse magnetic correlation length of La<sub>2</sub>CuO<sub>4</sub>. The solid line is Eq. (6) with  $J=135$  meV. The Néel and structural transition temperatures are indicated by arrows.

$$\frac{\xi}{a} = \frac{e}{8} \frac{c/a}{2\pi\rho_s} e^{2\pi\rho_s/T} \left[ 1 - \frac{1}{2} \left( \frac{T}{2\pi\rho_s} \right) + \mathcal{O} \left( \frac{T}{2\pi\rho_s} \right)^2 \right], \quad (6)$$

CHN: Phys. Rev. B 39, 2344, 1989

## Two-dimensional quantum Heisenberg antiferromagnet at low temperatures

Sudip Chakravarty\*

*Department of Physics, State University of New York at Stony Brook, Stony Brook, New York 11794*

Bertrand I. Halperin and David R. Nelson

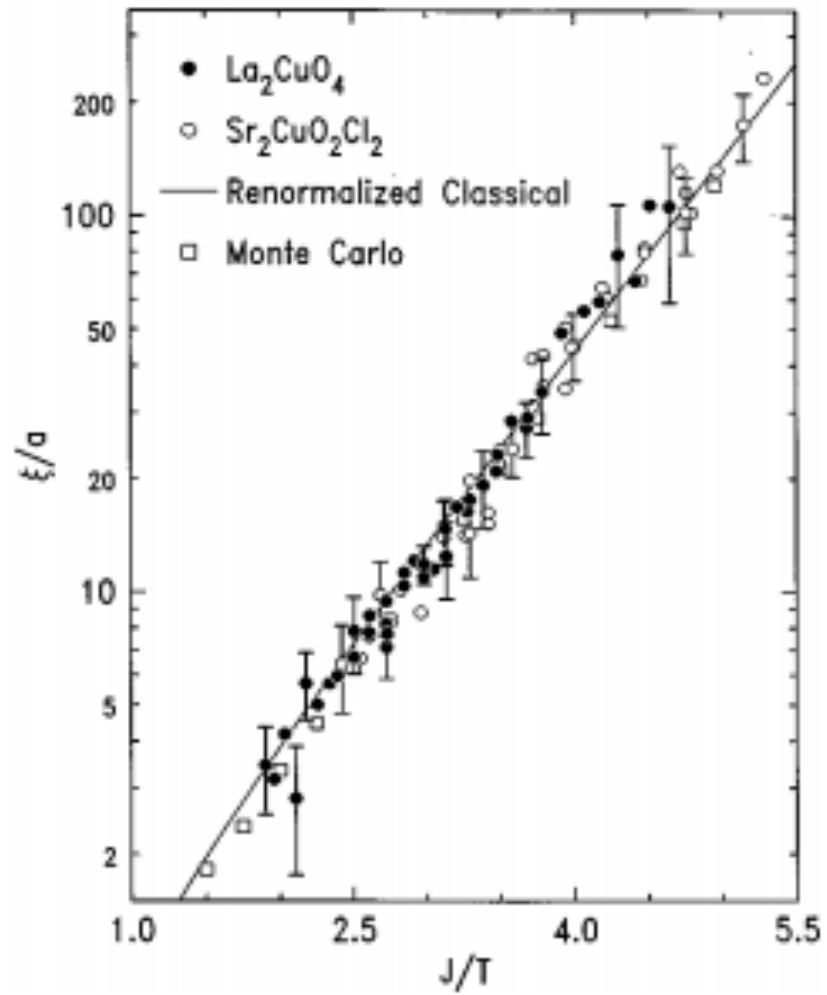
*Lyman Laboratory of Physics, Harvard University, Cambridge, Massachusetts 02138*

(Received 18 August 1988)

It is argued that the long-wavelength, low-temperature behavior of a two-dimensional quantum Heisenberg antiferromagnet can be described by a quantum nonlinear  $\sigma$  model in two space plus one time dimension, at least in the range of parameters where the model has long-range order at zero temperature. The properties of the quantum nonlinear  $\sigma$  model are analyzed approximately using the one-loop renormalization-group method. When the model has long-range order at  $T=0$ , the long-wavelength behavior at finite temperatures can be described by a purely classical model, with parameters renormalized by the quantum fluctuations. The low-temperature behavior of the correlation length and the static and dynamic staggered-spin-correlation functions for the quantum antiferromagnet can be predicted, in principle, with no adjustable parameters, from the results of simulations of the classical model on a lattice, combined with a two-loop renormalization-group analysis of the classical nonlinear  $\sigma$  model, a calculation of the *zero-temperature* spin-wave stiffness constant and uniform susceptibility of the quantum antiferromagnet, and a one-loop analysis of the conversion from a lattice cutoff to the wave-vector cutoff introduced by quantum mechanics when the spin-wave frequency exceeds  $T/\hbar$ . Applying this approach to the spin- $\frac{1}{2}$  Heisenberg model on a square lattice, with nearest-neighbor interactions only, we obtain a result for the correlation length which is in good agreement with the data of Endoh *et al.* on  $\text{La}_2\text{CuO}_4$ , if the spin-wave velocity is assumed to be  $0.67 \text{ eV } \text{\AA}/\hbar$ . We also argue that the data on  $\text{La}_2\text{CuO}_4$  cannot be easily explained by any model in which an isolated  $\text{CuO}_2$  layer would not have long-range antiferromagnetic order at  $T=0$ . Our theory also predicts a quasielastic peak of a few meV width at 300 K when  $k\xi \ll 1$  (where  $k$  is wave-vector transfer and  $\xi$  is the correlation length). The extent to which this dynamical prediction agrees with experiments remains to be seen. In an appendix, we discuss the effect of introducing a frustrating second-nearest-neighbor coupling for the antiferromagnet on the square lattice.



### Magnetic Correlation Length



# Rise of the iron pnictides

**J|A|C|S**  
COMMUNICATIONS

Published on Web 02/23/2008

**Iron-Based Layered Superconductor  $\text{La}[\text{O}_{1-x}\text{F}_x]\text{FeAs}$  ( $x = 0.05\text{--}0.12$ )  
with  $T_c = 26\text{ K}$**

Yoichi Kamihara,<sup>\*,†</sup> Takumi Watanabe,<sup>‡</sup> Masahiro Hirano,<sup>†,\$</sup> and Hideo Hosono<sup>†,‡,\$</sup>

# Rise of the iron pnictides

**J|A|C|S**  
COMMUNICATIONS

nature

Vol 453|15 April 2008|doi:10.1038/nature06972

LETTERS

2)

CHIN.PHYS.LETT.

Vol. 25, No. 6 (2008) 2215



A LETTERS JOURNAL EXPLORING  
THE FRONTIERS OF PHYSICS

September 2008

EPL, 83 (2008) 67006  
doi: 10.1209/0295-5075/83/67006

www.epjjournal.org

## Thorium-doping-induced superconductivity up to 56 K in $Gd_{1-x}Th_xFeAsO$

CAO WANG, LINJUN LI, SHUN CHI, ZENGWEI ZHU, ZHI REN, YUKE LI, YUETAO WANG, XIAO LIN,  
YONGKANG LUO, SHUAI JIANG, XIANGFAN XU, GUANGHAN CAO<sup>(a)</sup> and ZHU'AN XU<sup>(b)</sup>

**Berkeley**  
UNIVERSITY OF CALIFORNIA

# Rise of the iron pnictides

**J|A|C|S**  
COMMUNICATIONS

nature

Vol 453|15 April 2008|doi:10.1038/nature06972

LETTERS

2)

CHIN.PHYS.LETT.

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THE FRONTIERS OF PHYSICS

September 2008

EPL, 83 (2008) 67006

[www.epljournal.org](http://www.epljournal.org)

doi: 10.1209/0295-5075/83/67006

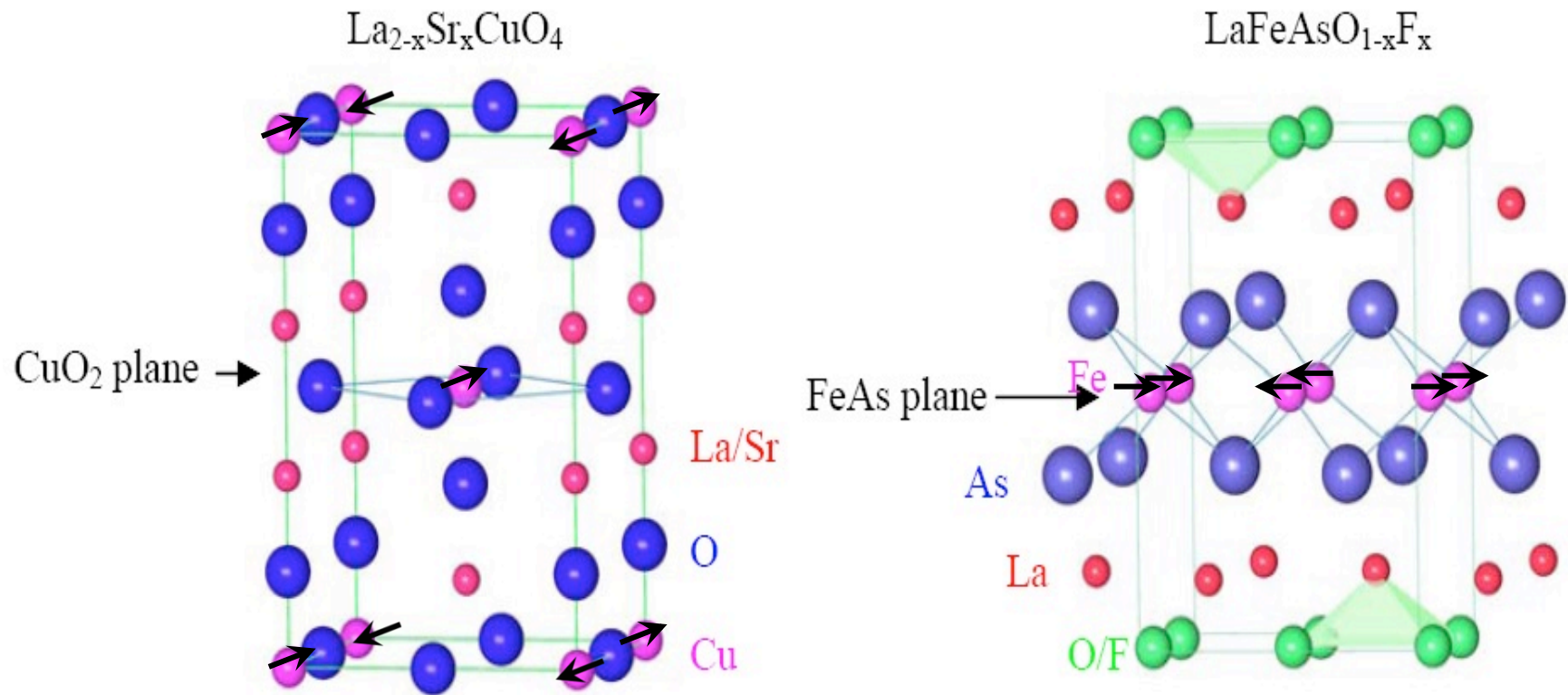
## **Superconductivity in single-layer films of FeSe with a transition temperature above 100 K**

Jian-Feng Ge<sup>1</sup>, Zhi-Long Liu<sup>1</sup>, Canhua Liu<sup>1\*</sup>, Chun-Lei Gao<sup>1</sup>, Dong Qian<sup>1</sup>, Qi-Kun Xue<sup>2\*</sup>, Ying  
Liu<sup>1,3</sup>, Jin-Feng Jia<sup>1\*</sup>

arxiv:1406.3435

**Berkeley**  
UNIVERSITY OF CALIFORNIA

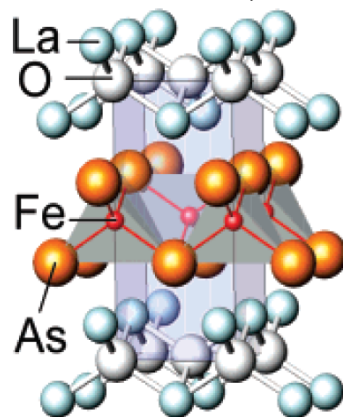
# Common features of high $T_c$ superconductors



# New Iron Based Superconductors

ZrCuSiAs-type

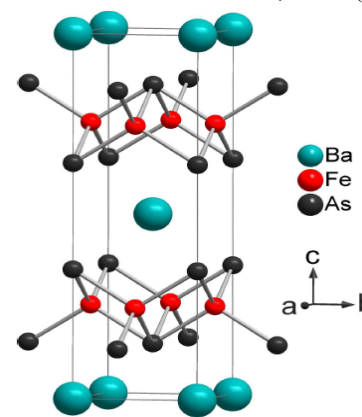
Kamihara et al., JACS (2008)



$(\text{RE})\text{FeAsO}_{1-x}\text{F}_x$  ( $T_c \sim 55\text{K}$ )<sub>max</sub>

ThCr<sub>2</sub>Si<sub>2</sub>-type

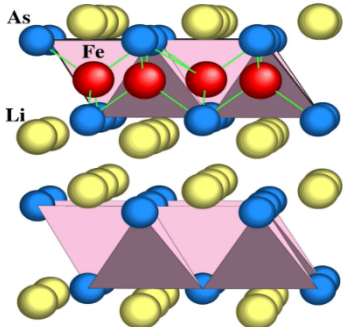
Rotter et al., PRL (2008)



$\text{Ba}_{1-x}\text{K}_x\text{Fe}_2\text{As}_2$  ( $T_c \sim 38\text{K}$ )<sub>max</sub>

Cu<sub>2</sub>Sb-type

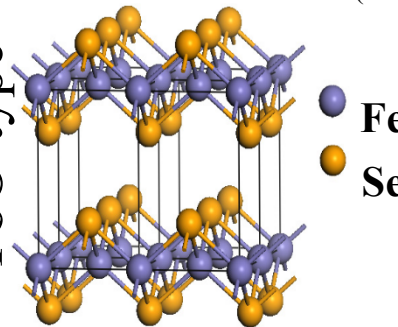
Nekrasov et al. JETP (2008)



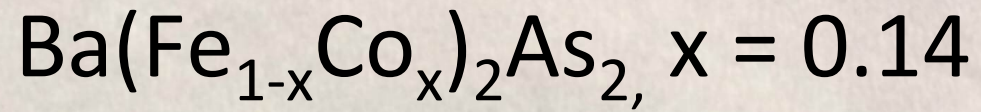
$\text{LiFeAs}$  ( $T_c \sim 18\text{K}$ )<sub>max</sub>

PbO-type

Hsu et al. Arxiv (2008)



$\text{FeSe}_{1-x}\text{Te}$  ( $T_c \sim 15\text{K}$ )<sub>max</sub>



0.46 g



0.96 g



0.55 g



70

80

90

100

110

120



0.53 g



0.35 g



0.1 g

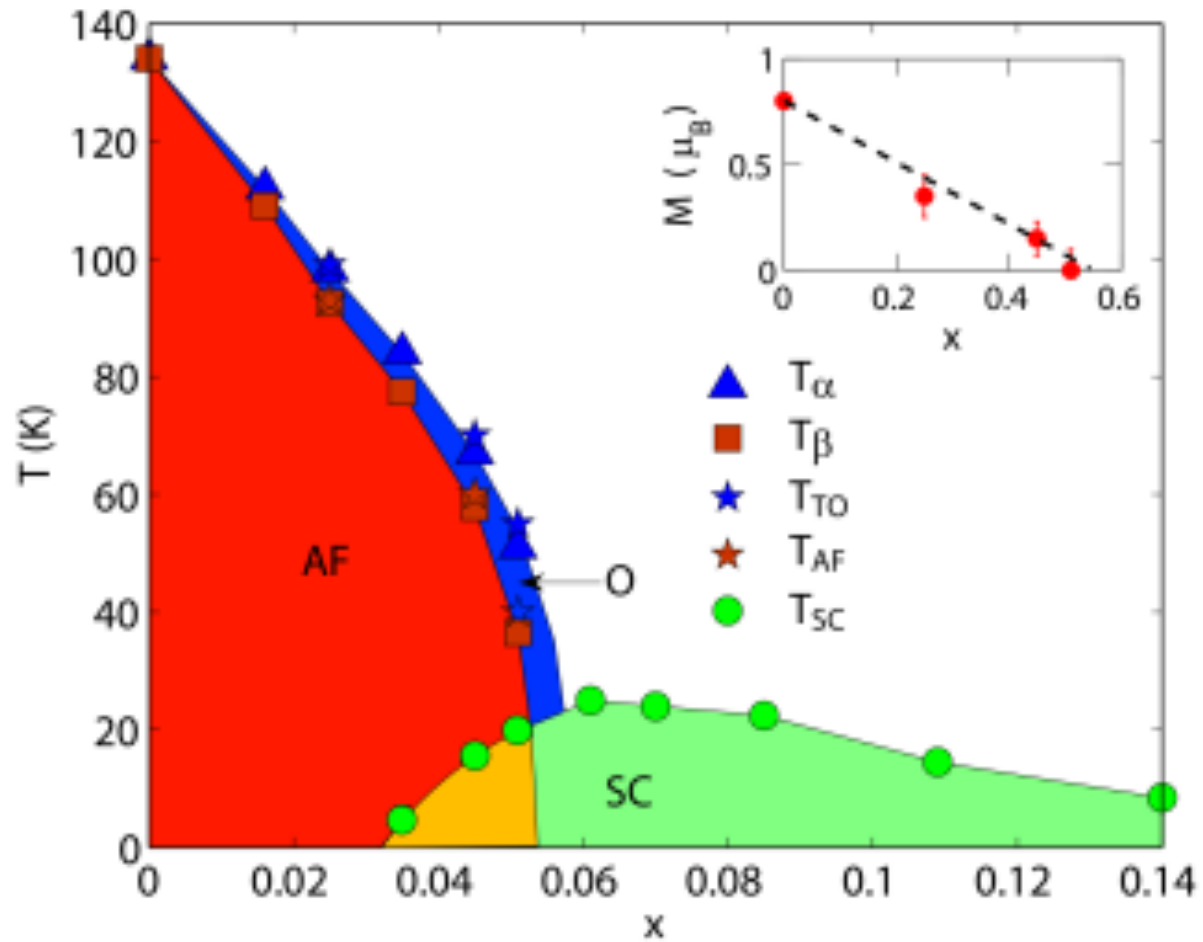
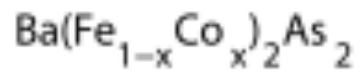
$\text{Rb}_{0.8}\text{Fe}_{1.5}\text{SeS}$

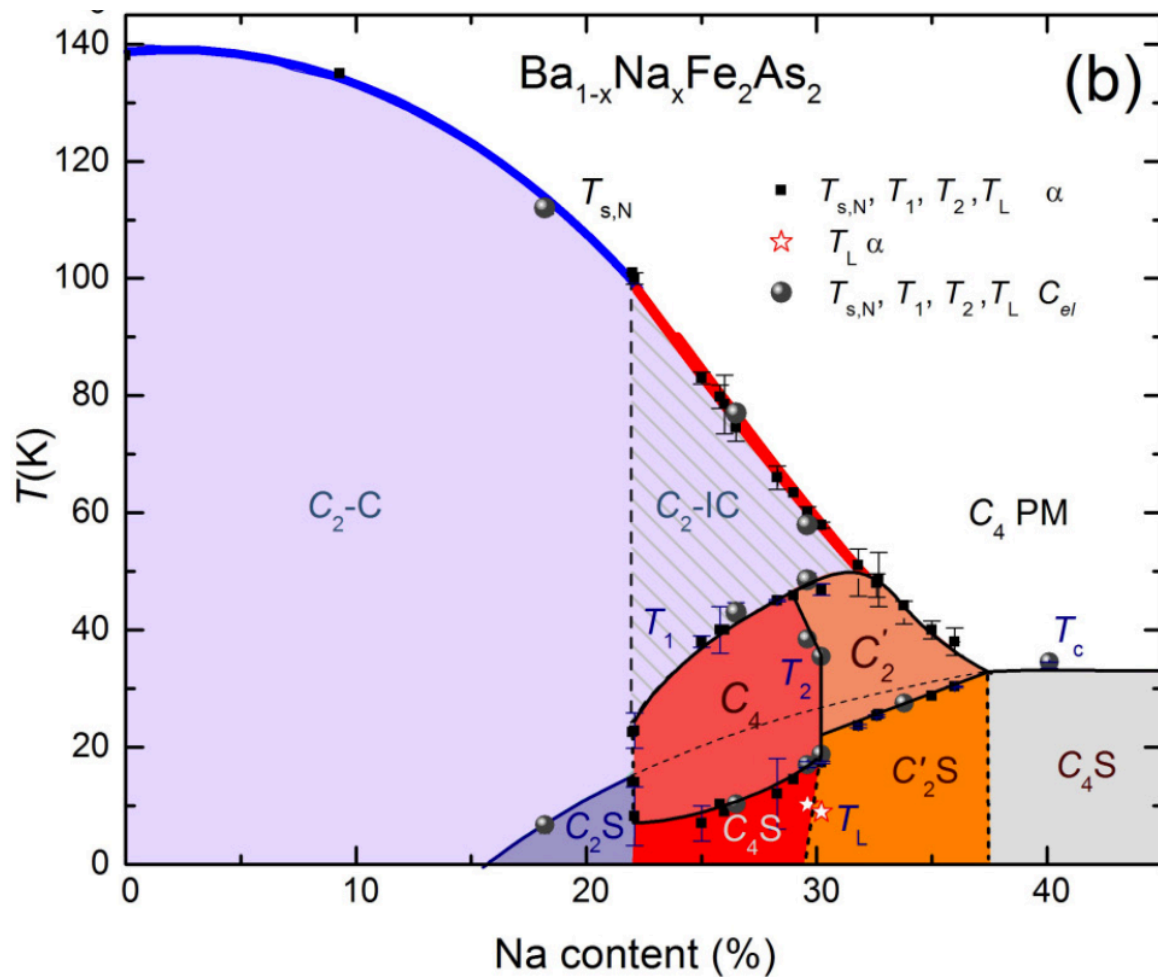
3 grams

Grown with Bridgman furnace







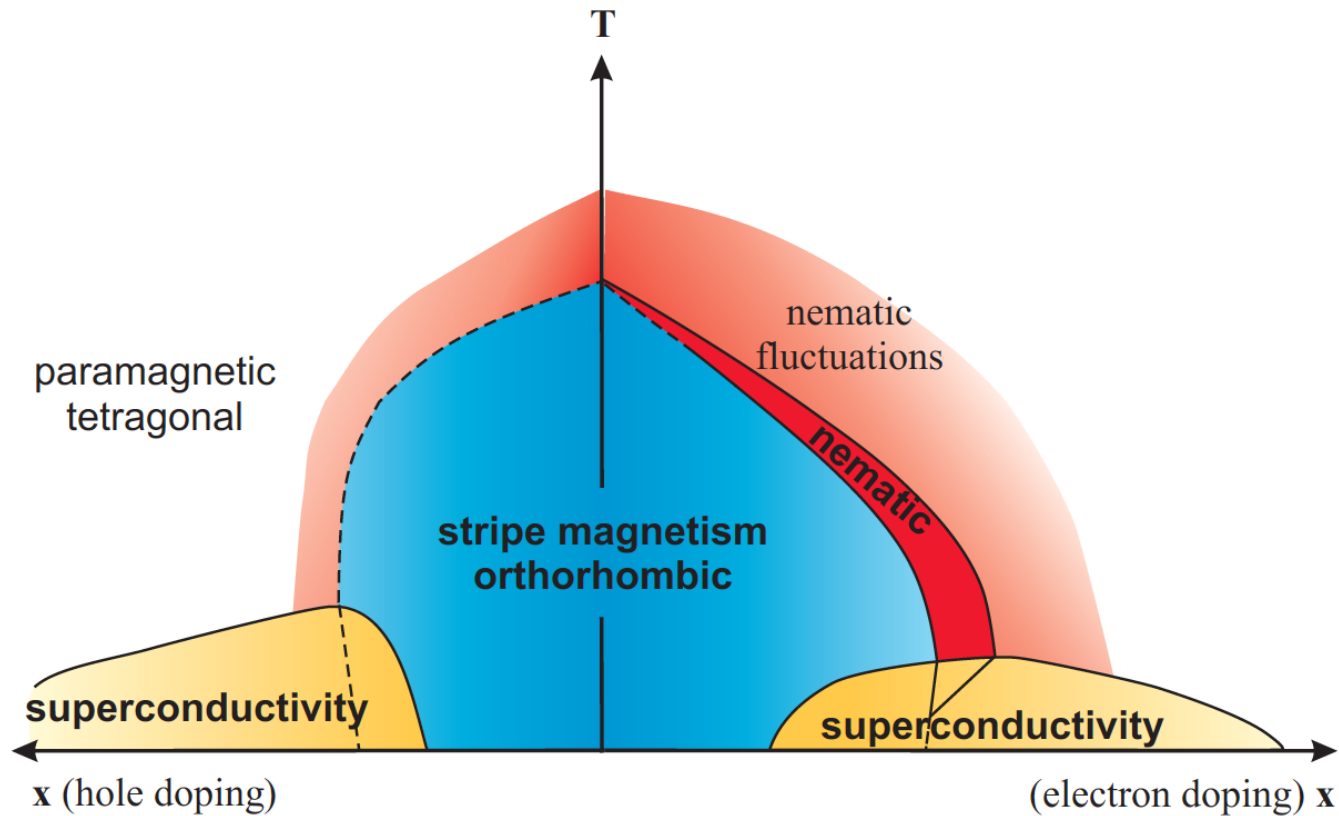


A. Bohmer *et al.* Nat. Comm. 6, 7911 (2015)

L. Wang *et al.* Phys. Rev. B 93, 014514 (2016)

# Schematic phase diagram of the iron pnictides

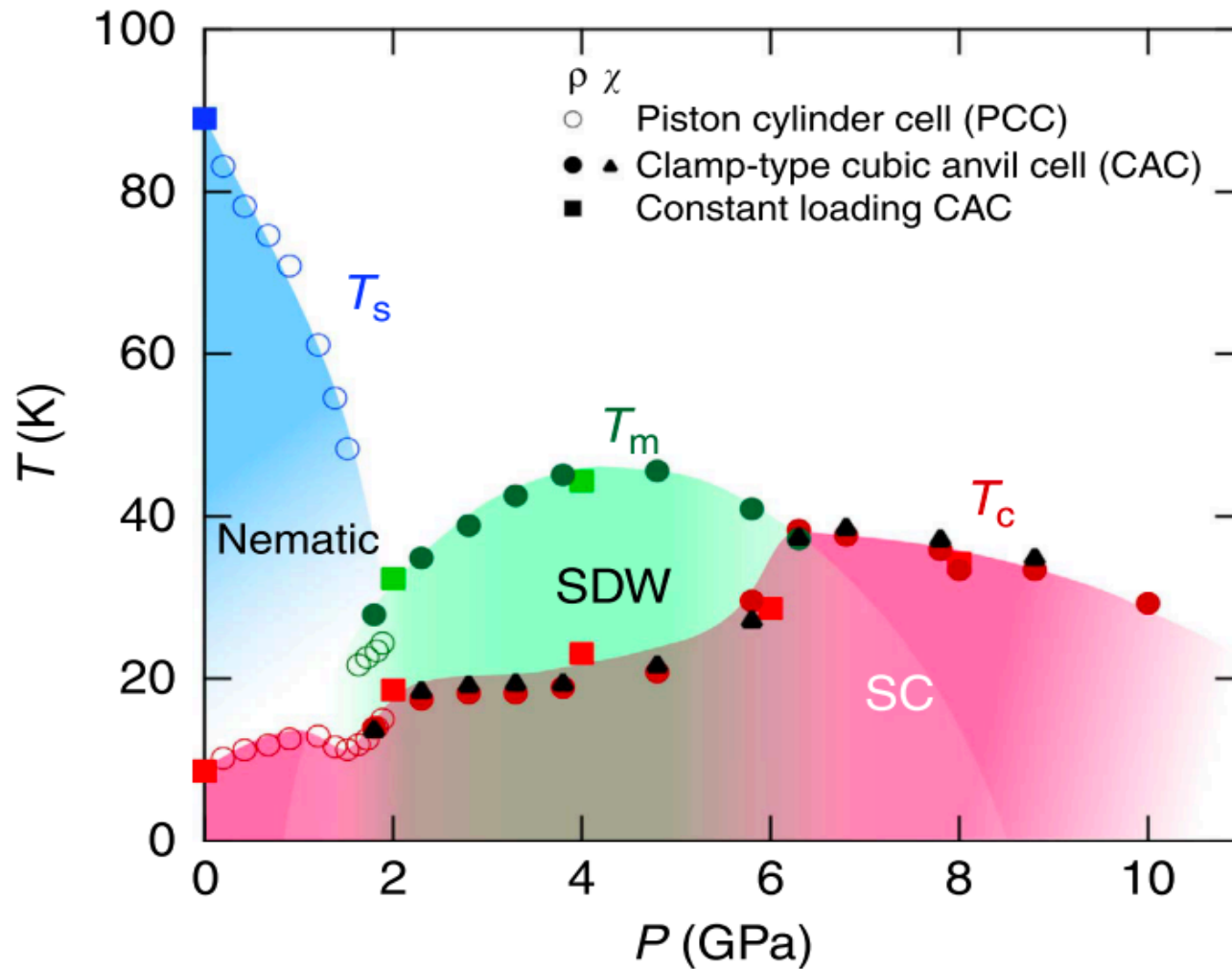
- what is the origin of the nematic instability?
- what is the importance of nematicity to superconductivity?



***nematic transition = tetragonal-to-orthorhombic transition***

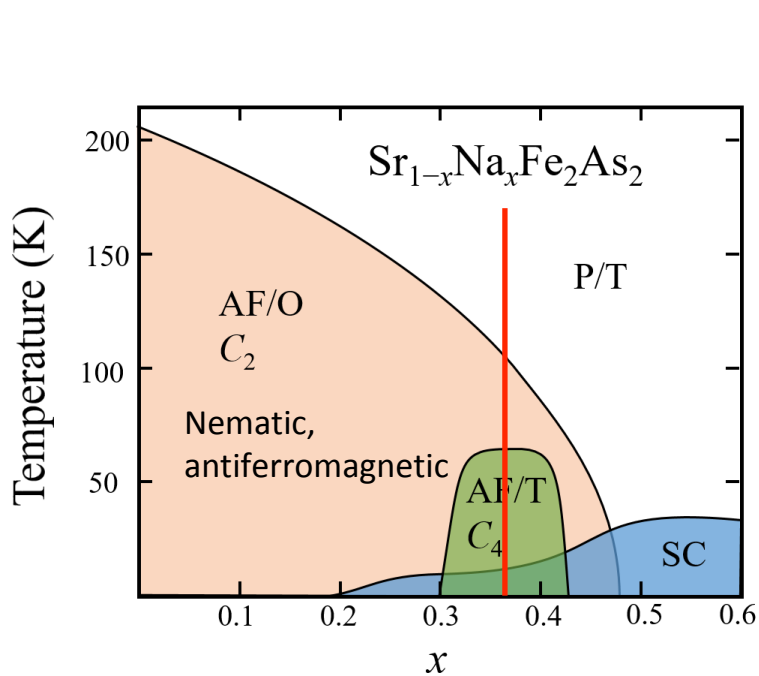
Rafael Fernandes

# FeSe phase diagram: temperature vs pressure

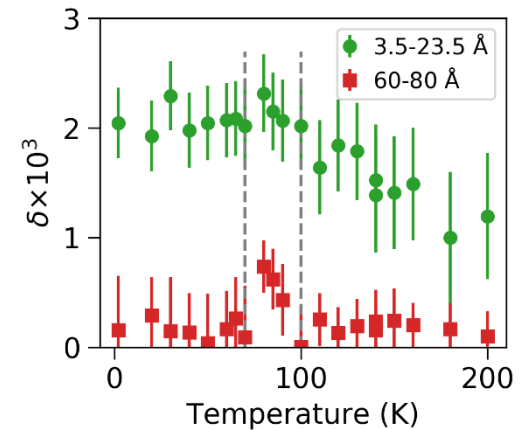
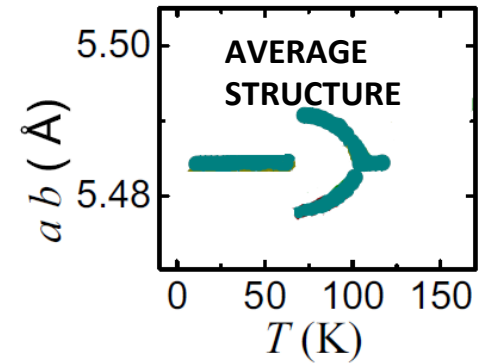
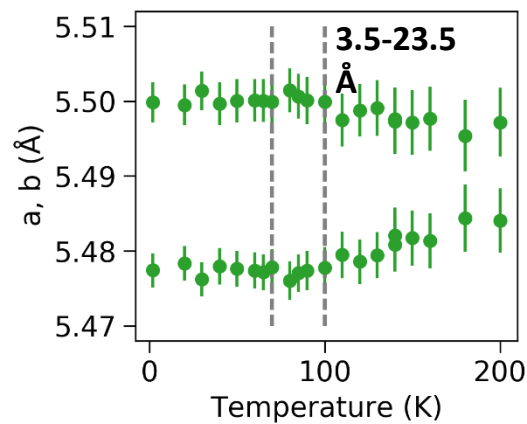
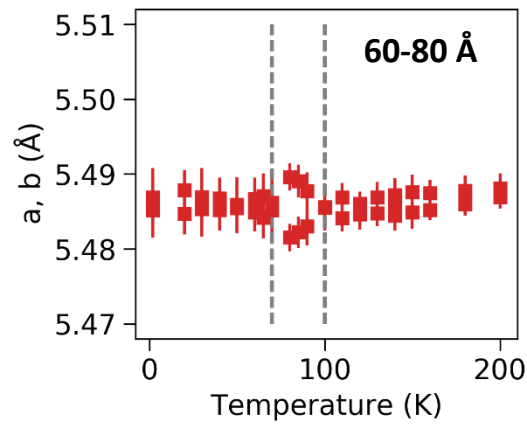


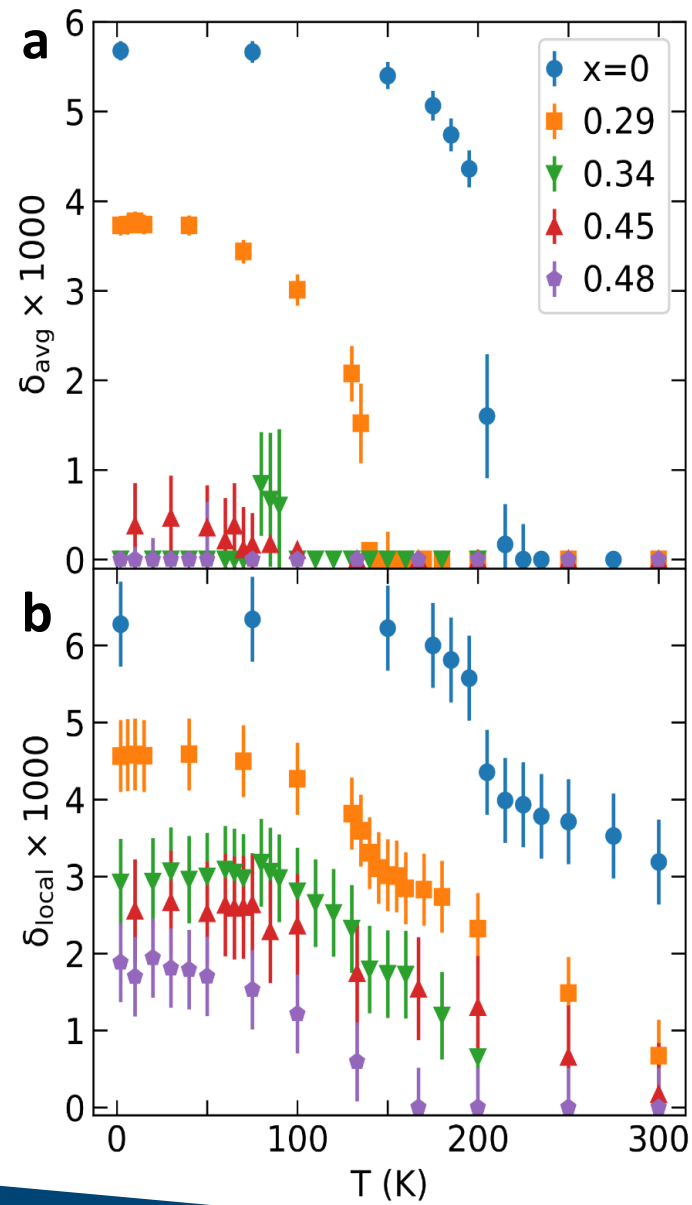
J. Sun *et al.* Nat. Comm. 7 12146 (2016)

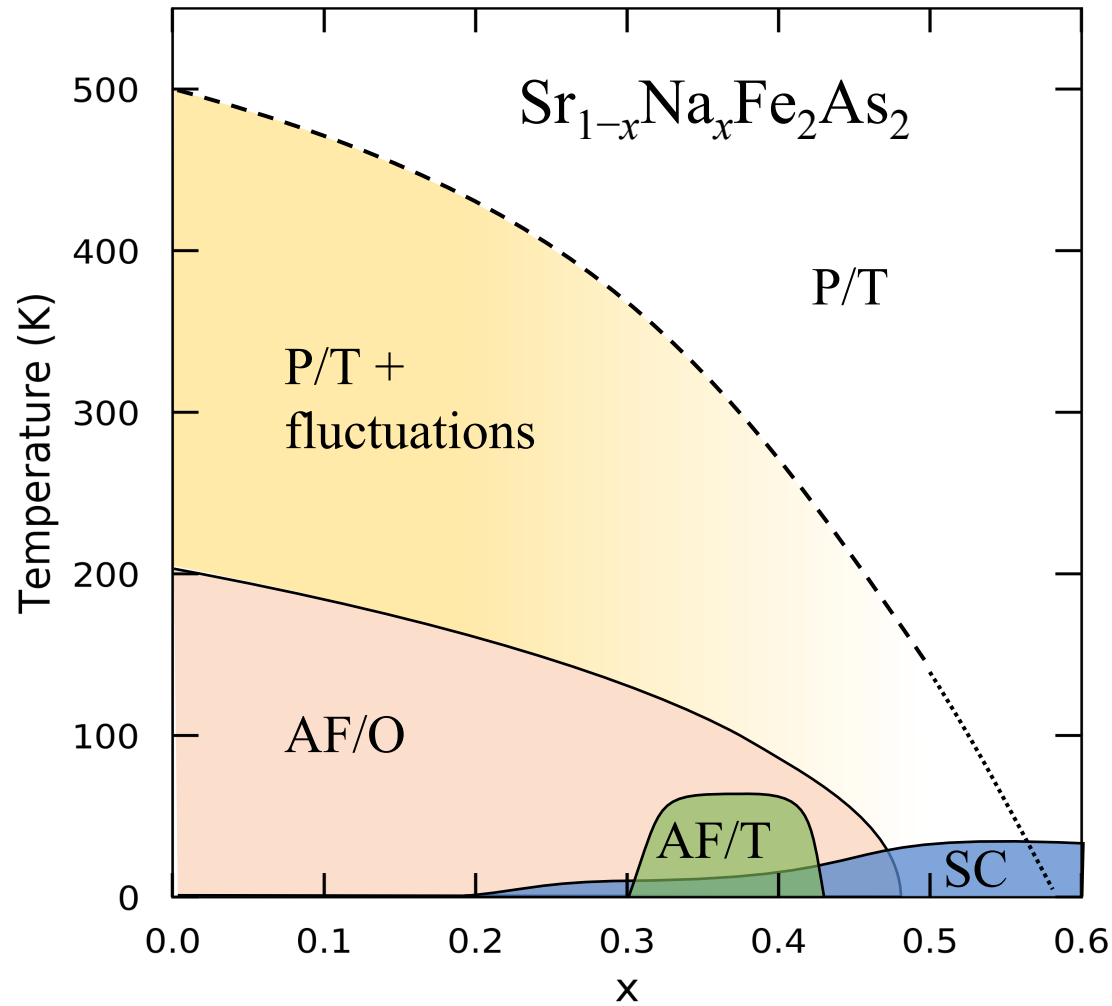
# Structure of $\text{Sr}_{1-x}\text{Na}_x\text{Fe}_2\text{As}_2$ : $x = 0.34$ (with magnetic $C_4$ phase)



Frandsen et al, PRL **119**, 187001 (2017)



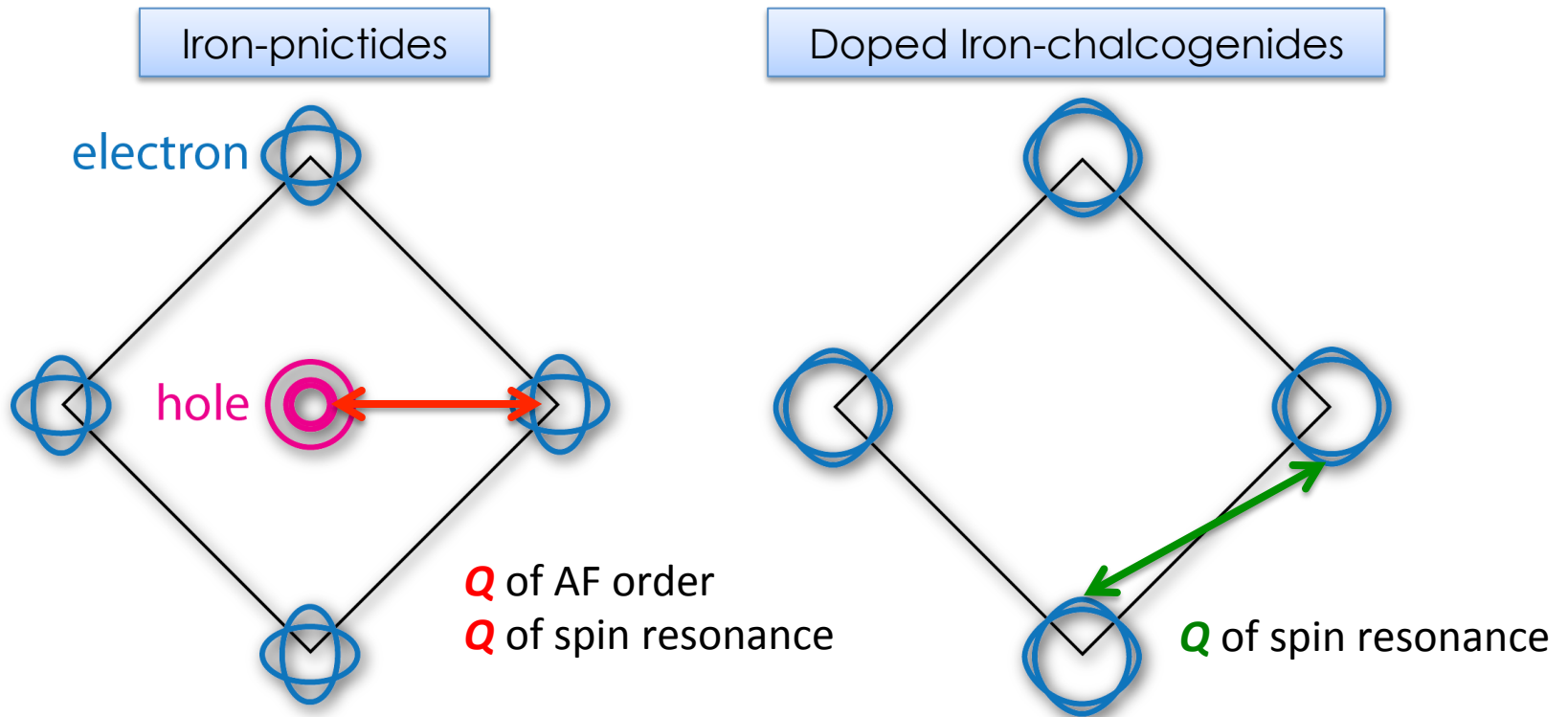




Note:  
Disorder acts  
as a random  
nematic field

→ The physics  
of the 3D  
random field  
Ising Model

# The curious case of doped iron chalcogenide superconductors

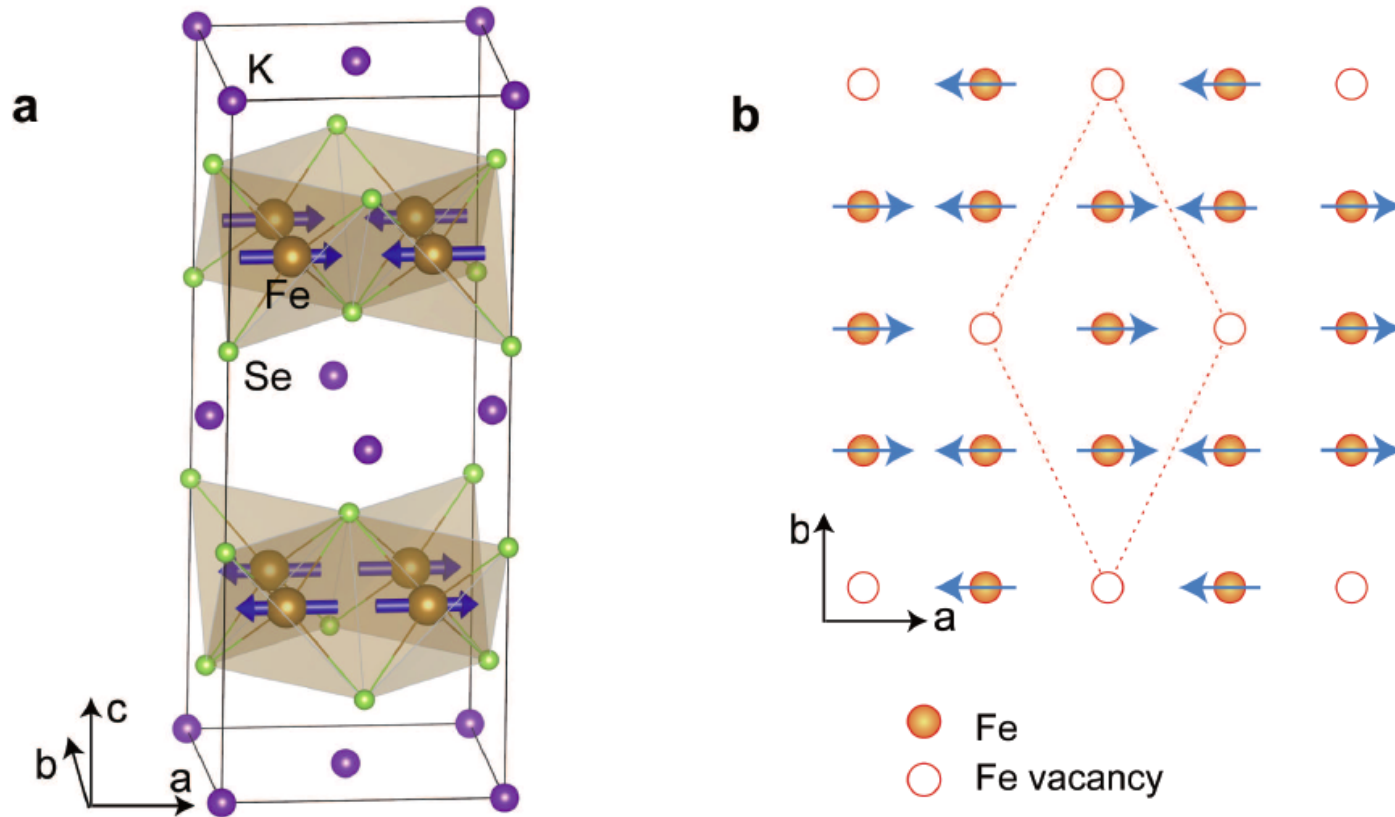


Spin fluctuations at the Fermi surface nesting vector  $(\pi, \pi)$

Lack of electron and hole Fermi surface nesting conditions



# Stripe type AF order and rhombus iron vacancy order of semiconducting $K_{0.85}Fe_{1.54}Se_2$



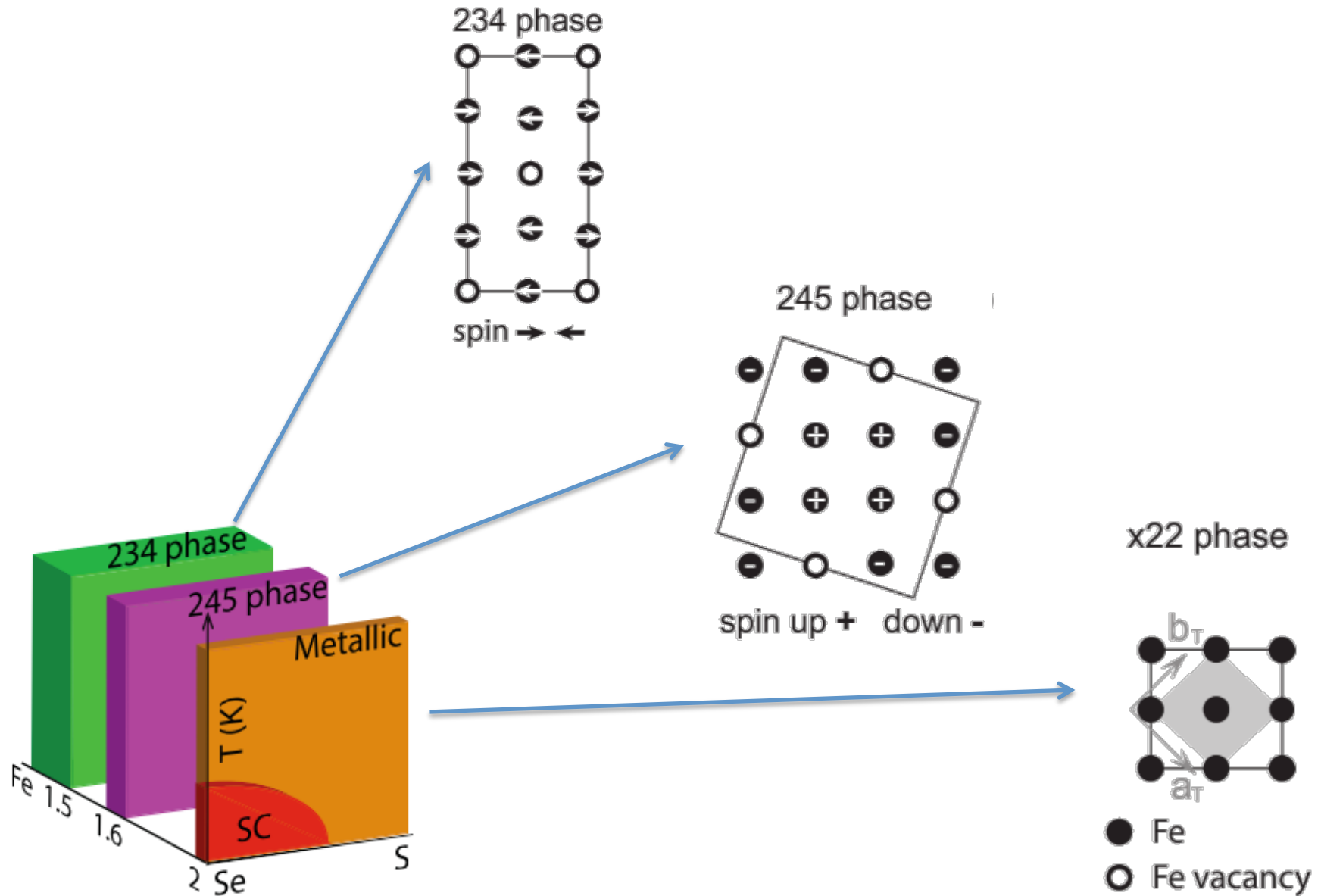
J. Zhao et al. arXiv:1205.5992 (2012)

TABLE I.

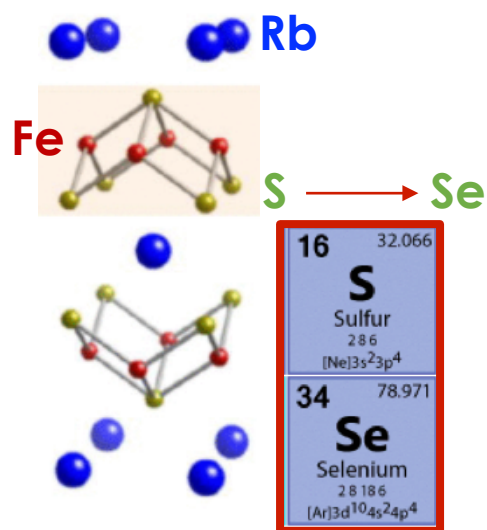
The magnetic exchange couplings and spin states in the stripe AF order of iron pnictides and chalcogenides [13,16,27].

Compounds	$SJ_{1a}$	$SJ_{1b}$	$SJ_2$ (meV)	$S$	$M$ ( $\mu_B$ )	$T_N$ (K)
CaFe <sub>2</sub> As <sub>2</sub>	50(10)	-6(5)	19(4)	1/2	0.80	173
BaFe <sub>2</sub> As <sub>2</sub>	59(2)	-9(2)	14(1)	1/2	0.87	143
SrFe <sub>2</sub> As <sub>2</sub> (L)	31(1)	-5(5)	22(1)	0.30	0.94	220
SrFe <sub>2</sub> As <sub>2</sub> (H)	39(2)	-5(5)	27(1)	0.69	0.94	220
K <sub>0.85</sub> Fe <sub>1.54</sub> Se <sub>2</sub>	38(7)	-11(5)	19(2)		2.8	280
Rb <sub>0.8</sub> Fe <sub>1.5</sub> S <sub>2</sub>	42(5)	-20(2)	17(2)	2	2.8(0.5)	265

# Different phases in $\text{Rb}_x\text{Fe}_y\text{Se}_{2-z}\text{S}_z$

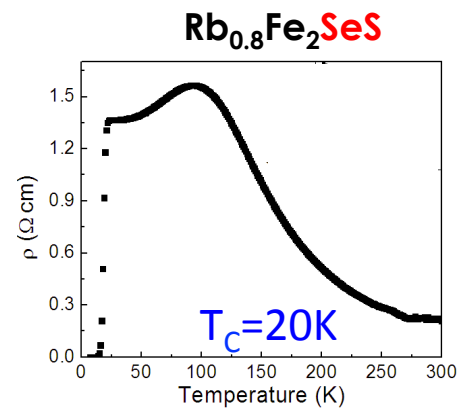
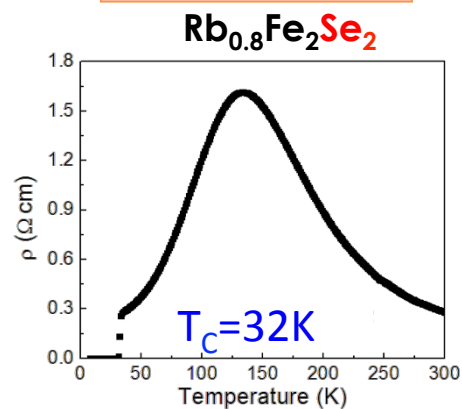


# Correlation-tuned superconductivity in $\text{Rb}_x\text{Fe}_2(\text{Se},\text{S})_2$

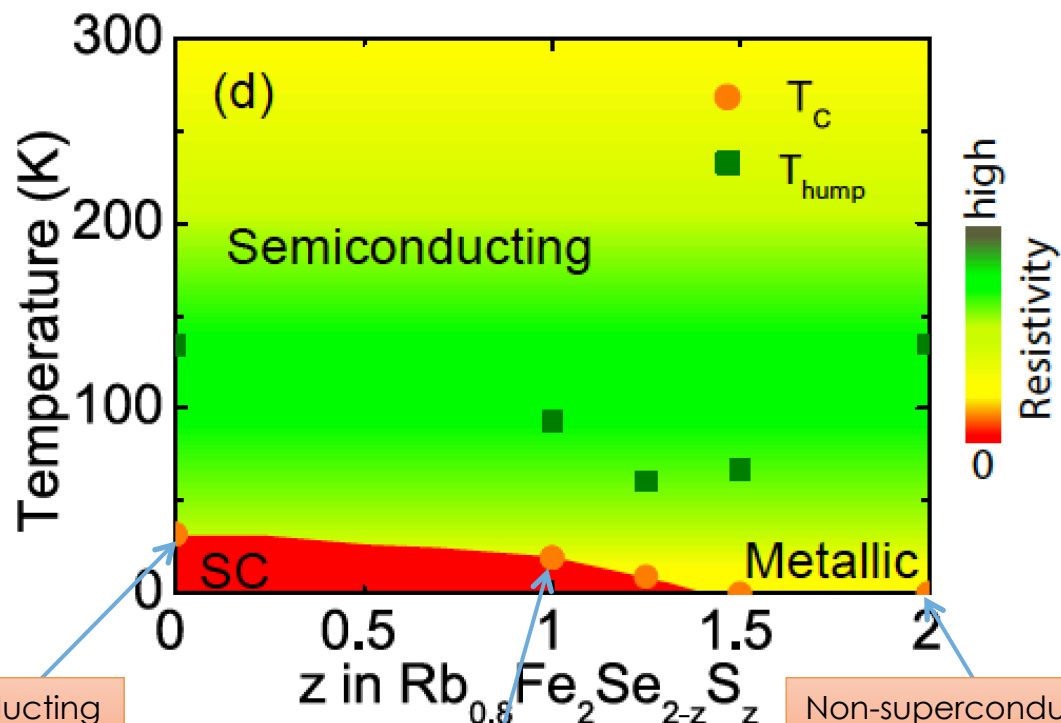
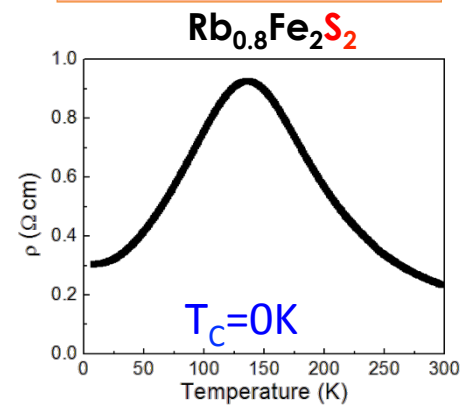


No charge carrier doping

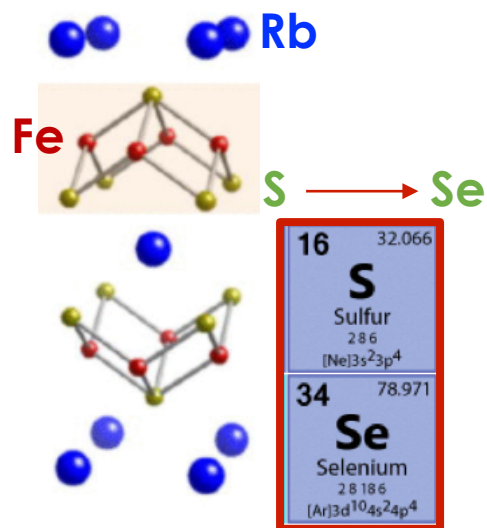
Superconducting



Non-superconducting

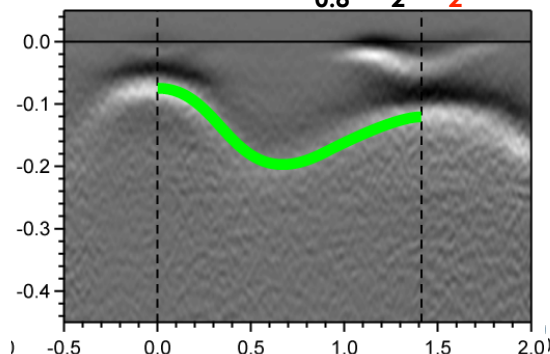


# Correlation-tuned superconductivity in $\text{Rb}_x\text{Fe}_2(\text{Se},\text{S})_2$



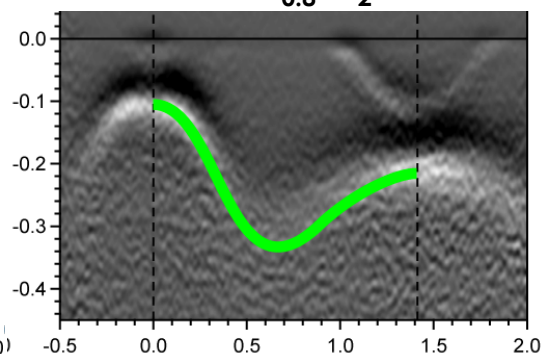
No charge carrier doping

Superconducting



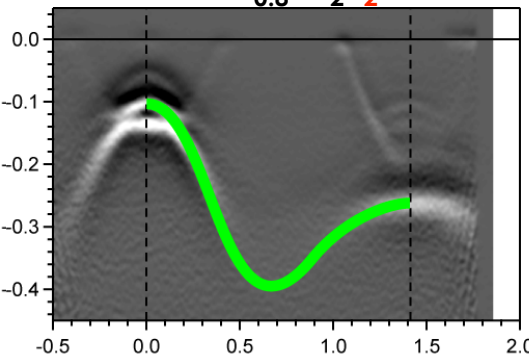
slower electrons  
– strong correlations

Non-superconducting

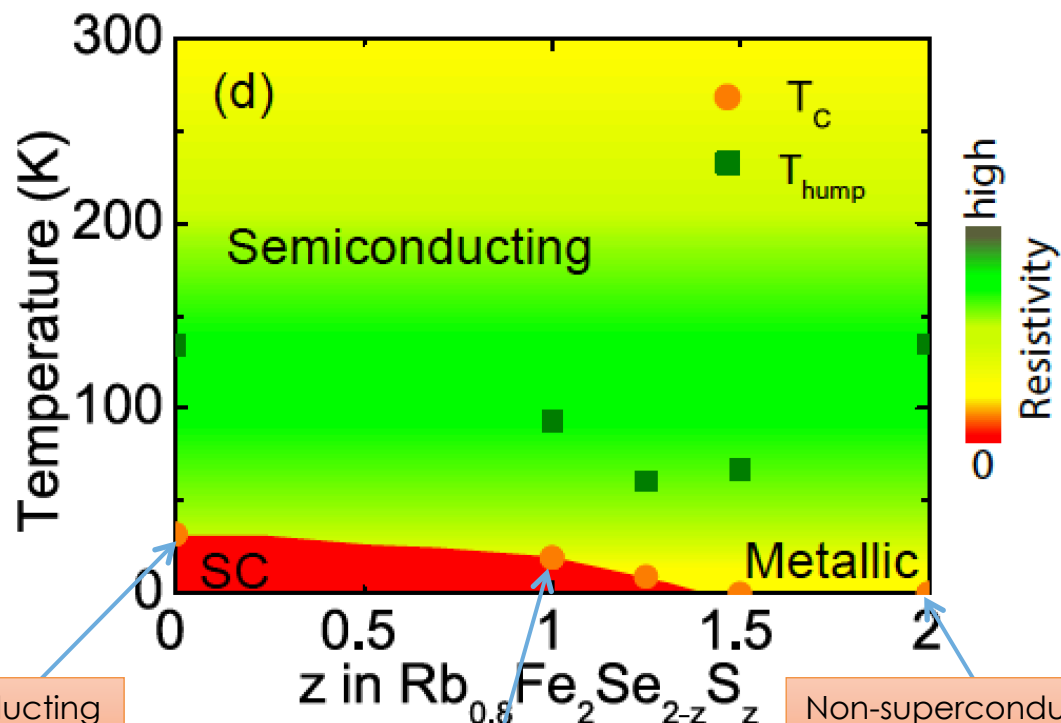


band slope ~ electron velocity

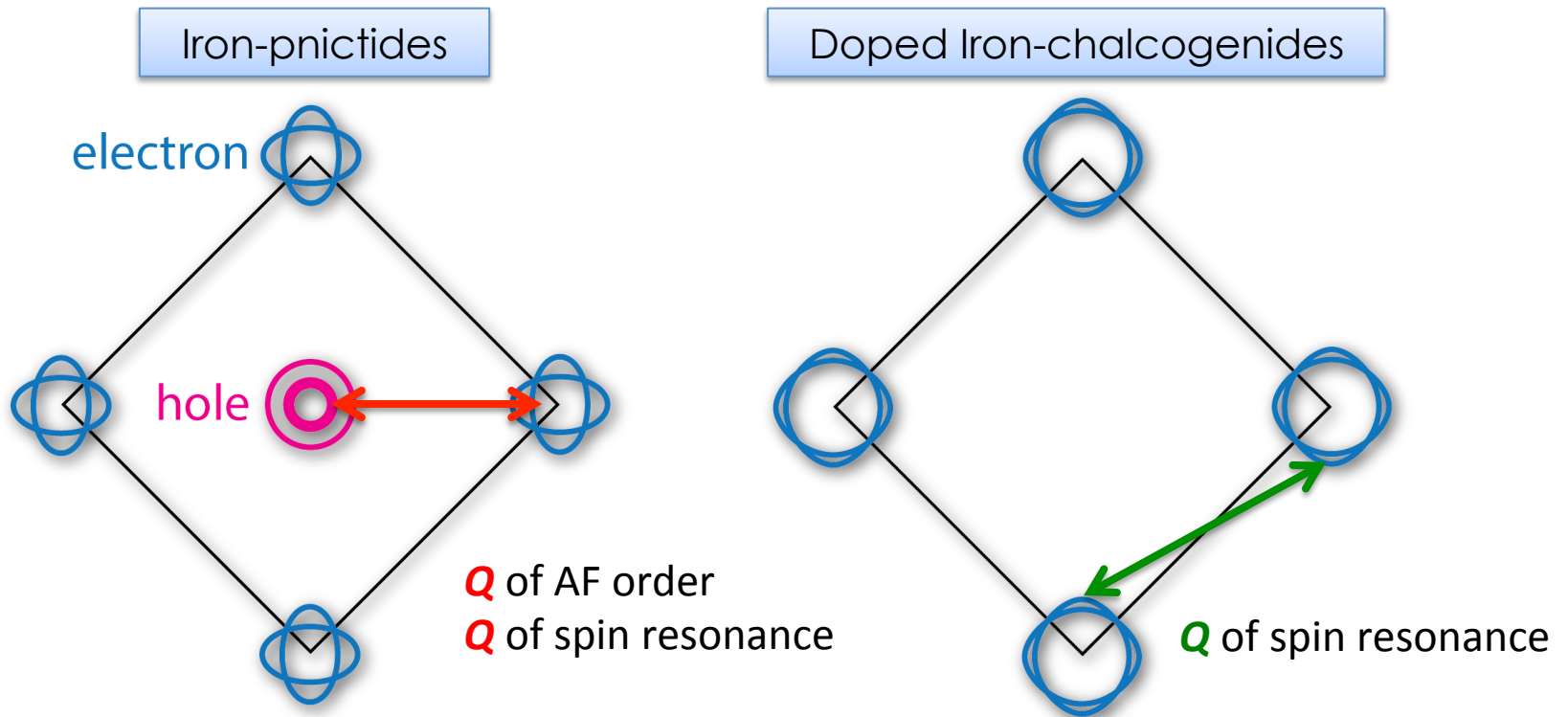
Non-superconducting



faster electrons  
– weak correlations



# The curious case of doped iron chalcogenide superconductors



Spin fluctuations at the Fermi surface nesting vector  $(\pi, \pi)$

Lack of electron and hole Fermi surface nesting conditions

# Spin resonant mode in $K_xFe_{2-y}Se_2$

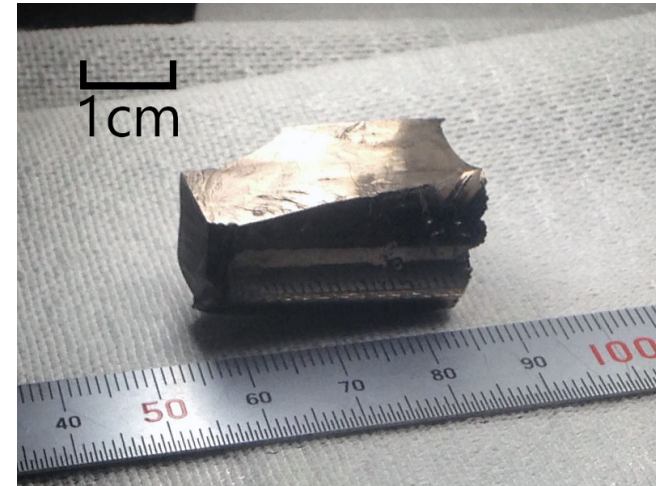
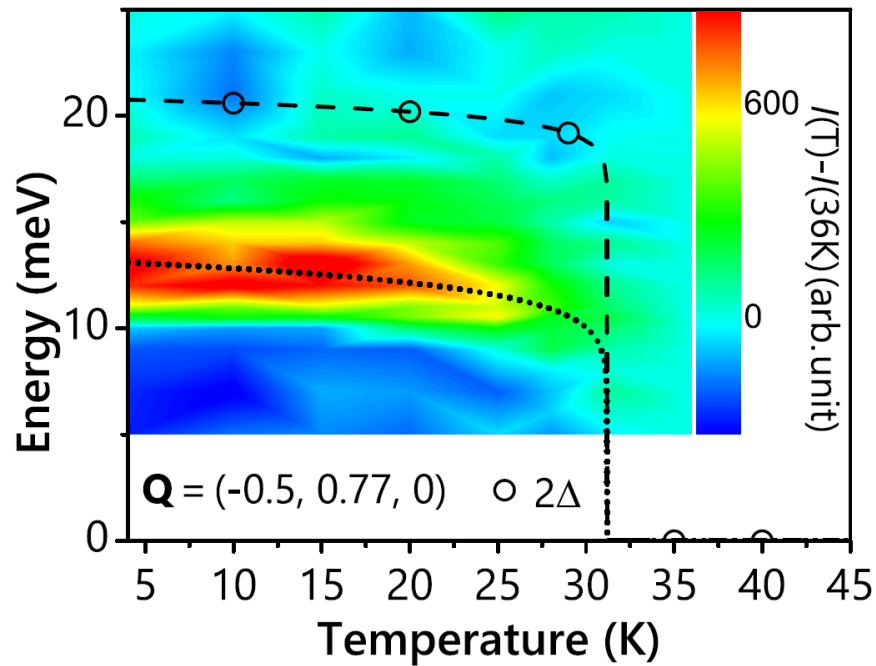


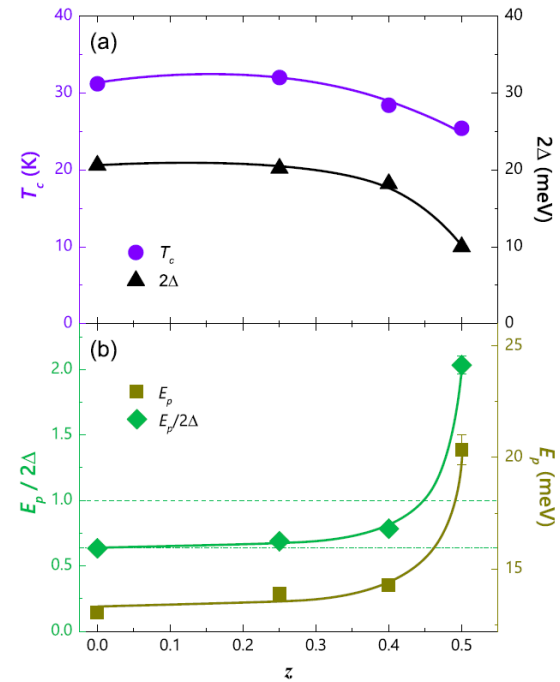
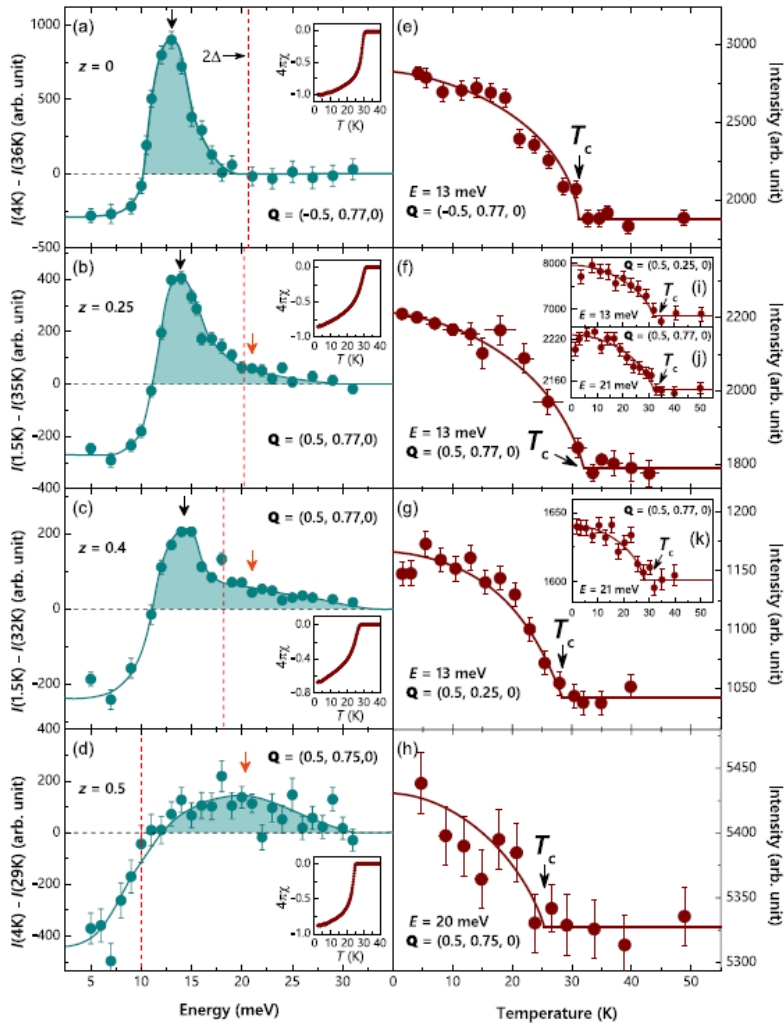
Photo of a  $KFe_{2-y}Se_2$  single crystal used for the INS measurement

Evidence for **sign-reversed** pairing symmetry in  $K_xFe_{2-y}Se_2$

Spin resonant mode below  $2\Delta$  in the superconducting state

Q. Wang *et al.*, *Phys. Rev. Lett.* **116**, 197004 (2016).

# Transition from sign-reversed to sign-preserved Cooper pairing in $K_xFe_{2-y}(Se_{1-z}S_z)_2$

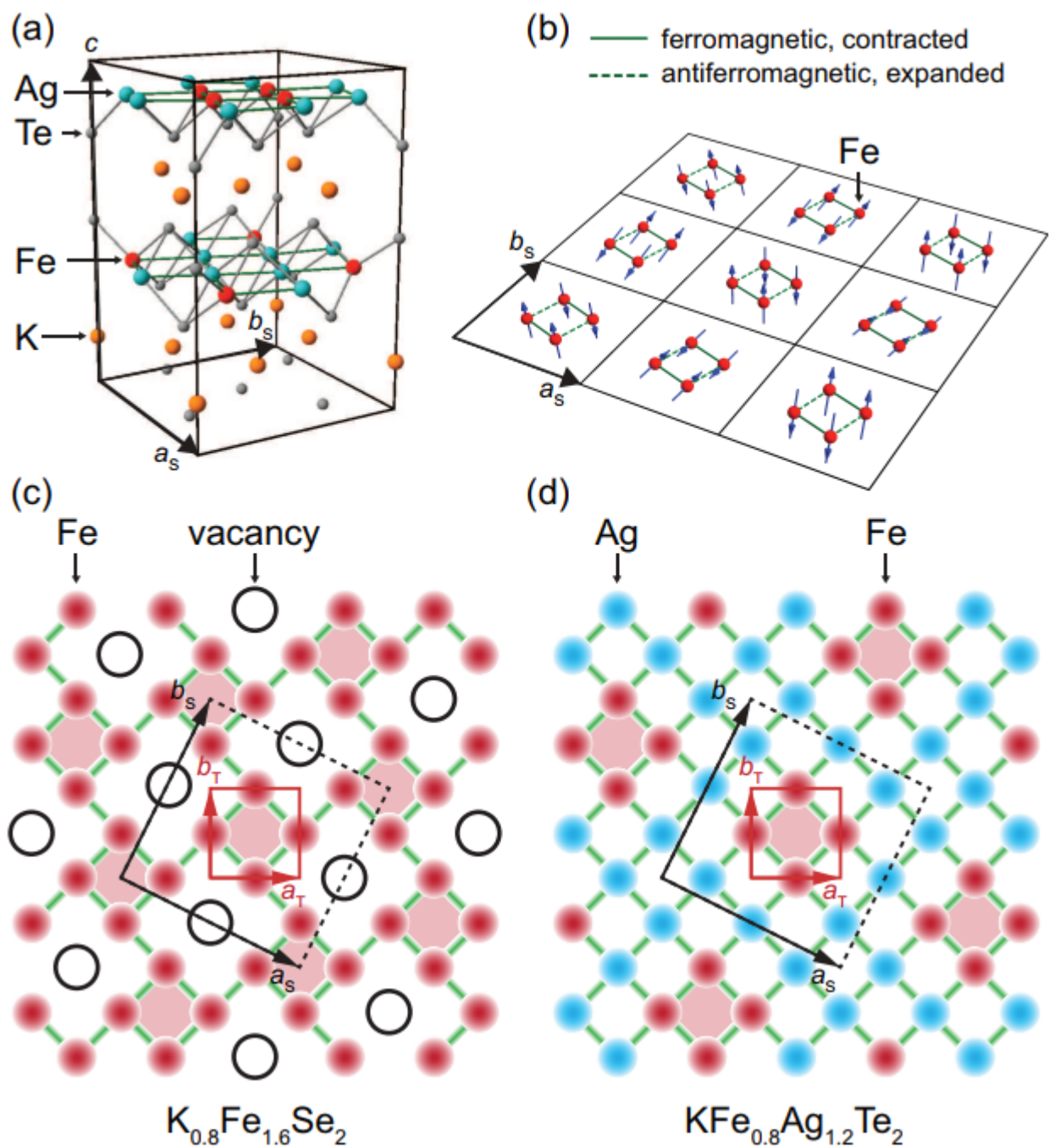


**sign-reversed ( $z=0$ ) to sign-preserved ( $z=0.5$ ) pairing**

Multiple pairing channels are required to understand the superconductivity in this system

Q. Wang *et al.*, *Phys. Rev. Lett.* **116**, 197004 (2016).





# Iron Pnictides and Chalcogenides

- Parent Materials are antiferromagnetic semi-metals or narrow band gap semiconductors.
- Magnetic and structural transitions are intimately connected. The structural (aka nematic) transition lifts the degeneracy of the  $dxz$  and  $dyz$  orbitals.
- Local orthorhombicity (a.k.a. nematicity) persists to temperatures far above the magnetic and superconducting transition temperatures.
- In the parent materials at low temperatures the magnetic exchange interactions appear to be universal.
- Antiferromagnetism and superconductivity compete.
- Phase diagrams have a remarkable similarity to those of the copper oxides.

Question: Is there a universal model?

# Lessons

- The most interesting developments in solid state physics, including especially superconductivity, have originated either from the discovery of new materials or the discovery of new properties of old materials.
- These materials require experimental studies with many different probes plus sophisticated theory. So far there has been no silver bullet.