

Dark Matter Searches and Fine-Tuning in Supersymmetry

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Laboratory for Elementary-Particle Physics

Bibliography

- Primary reference:

MP, Shakya, arXiv:1107.5048 [hep-ph]

- Precursors:

Mandic, Pierce, Gondolo, Murayama, hep-ph/0008022

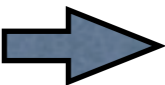
Kitano, Nomura, hep-ph/0606134

Feng, Sanford, arXiv:1009.3934 [hep-ph]

- Independent related recent work:

Amsel, Freese, Sandick, arXiv:1108.0448 [hep-ph]

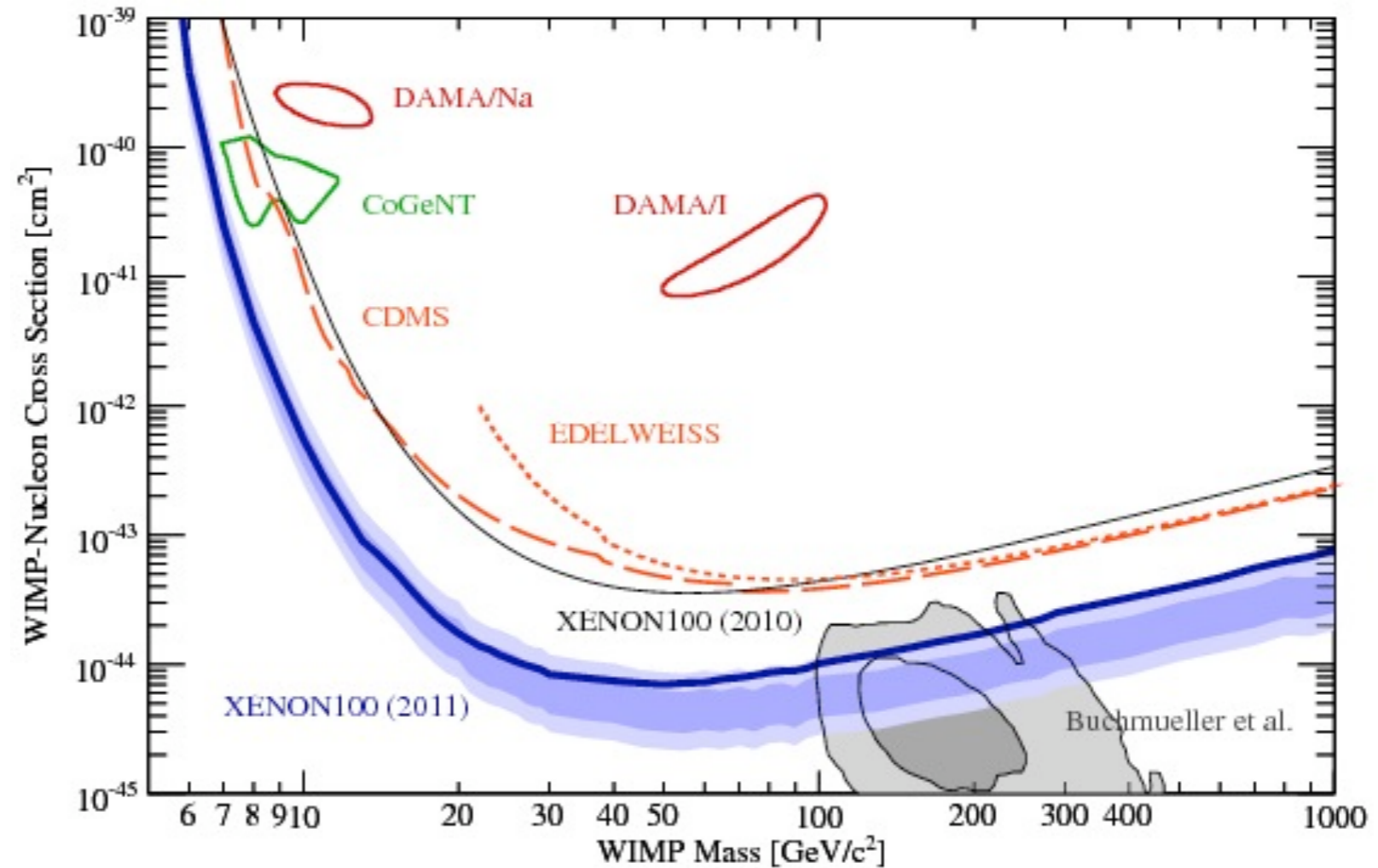
Introduction/Motivation

- Existence of dark matter **firmly established** through observing its gravitational interaction, on a variety of length scales from galaxy to Hubble
- No **SM** particle candidate for dark matter  solid observational evidence for **BSM physics!**
- Thermal relic with \sim weak-scale mass and \sim weak annihilation cross section (**WIMP**) naturally fits the observed dark matter abundance
- BS Models naturally contain WIMPs (e.g. SUSY with R-parity - stable **neutralino LSP**)

- Direct searches for WIMP dark matter have seen **significant improvements** recently, e.g. XENON100

[See e.g. B. Sadoulet's talk this morning]

“DD cross section” = spin-independent, elastic DM-proton scattering at zero mom. exchange

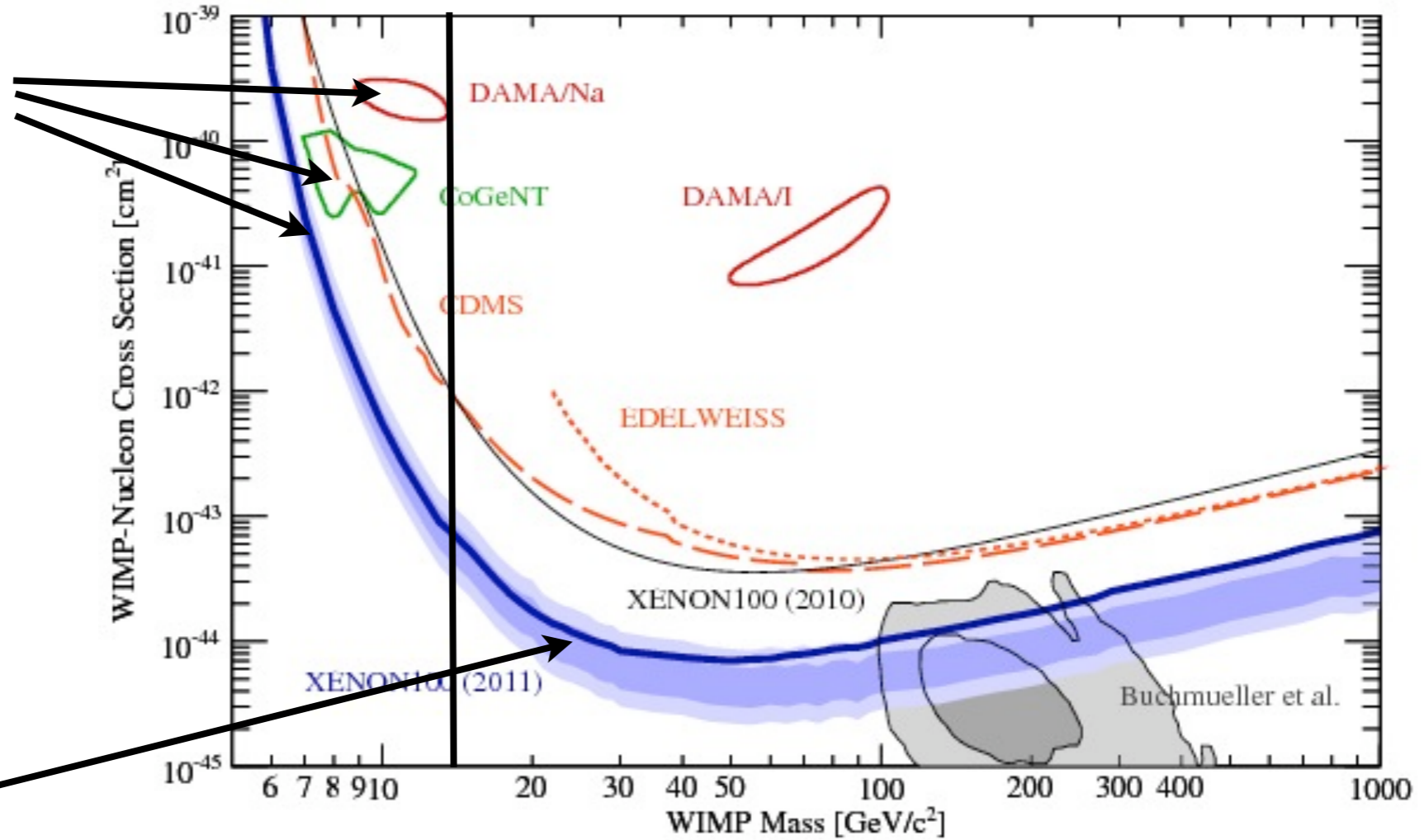


XENON-100, arXiv:1104.2549[astro-ph.CO]

- Direct searches for WIMP dark matter have seen significant improvements recently, e.g. XENON100

Confusion below 15 GeV:
contradictory (within
minimal WIMP framework)
claims

Clean exclusion
bound in 15 GeV -
1 TeV range

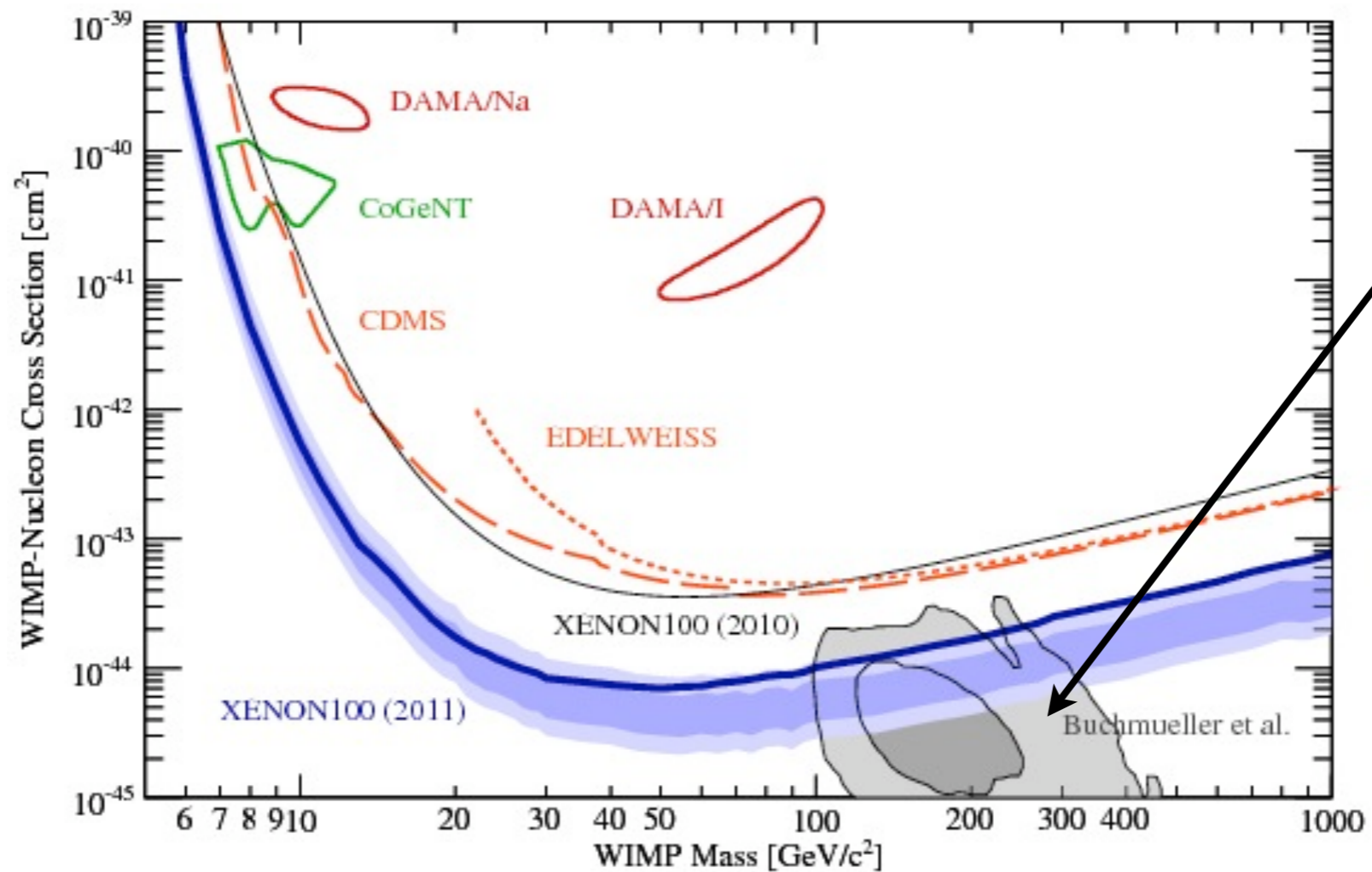


XENON-100, arXiv:1104.2549[astro-ph.CO]

MSSM Predictions for Direct Detection Cross Section

- Generic **problem** with MSSM predictions: large number (~ 100) of **free parameters**, allowed ranges are **highly uncertain** (esp. upper bounds on soft masses?)
- Standard approach: reduce parameter space by **assuming** high-scale unification, specific SUSY breaking model, etc.
- Most studied example, **mSUGRA**, has 5 parameters (and very serious issues with FCNC constraints!)

- Even then, the predictions span several orders of magnitude



“MSSM
predictions”
(mSUGRA
scatter plot)

XENON-100, arXiv:1104.2549[astro-ph.CO]

Our Approach

- Work with the MSSM defined in terms of **weak-scale** parameters, treat all of them as **independent** (as in pMSSM, “SUSY without prejudice”, etc.)
- Reduce # of parameters by assuming **absence of accidental cancellations** in the DD cross section
- Look for **correlations** between DD cross section and other physical quantities
- **Main question:** What’s special about points with low DD cross section?

Define “Accidental”?

- Intuitive definition:

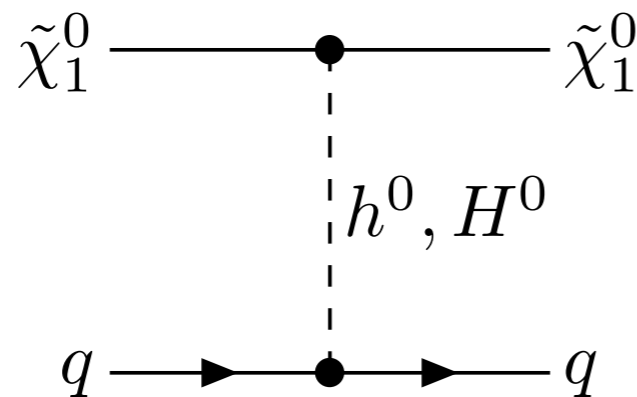
$$\sigma = c|\mathcal{M}_1 + \mathcal{M}_2|^2, \quad \mathcal{M}_1 \approx -\mathcal{M}_2, \quad \sigma \ll c|\mathcal{M}_1|^2$$

- More general definition:

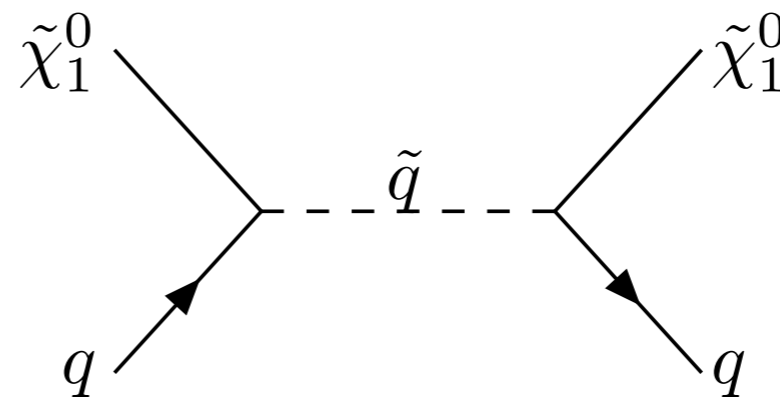
$$\sigma = \sigma(\{p_i\}), \quad \left| \frac{\partial \log \sigma}{\partial \log p_i} \right| \gg 1$$

- **Disclaimer:** this definition depends on which parameters are treated as “input”. A specific model of SUSY breaking may in fact predict relations that seem accidental from our (low-energy) point of view. It is however VERY hard to imagine this for the particular examples where we use this rule!

Direct Detection in the MSSM



(a)



(b)

$(M_1, M_2, \mu, \tan \beta, \underline{m_A})$

$(M_1, M_2, \mu, \tan \beta, \underline{m_{\tilde{q}_L}^2, m_{\tilde{q}_R}^2, A_q})$

With no accidental cancellations, sufficient to consider only one diagram class. **Choose (a).**

EWSB Fine-Tuning in the MSSM

- **Z mass** at the tree level in the MSSM:

$$m_Z^2 = -m_u^2 \left(1 - \frac{1}{\cos 2\beta}\right) - m_d^2 \left(1 + \frac{1}{\cos 2\beta}\right) - 2|\mu|^2$$

- Unless all terms on the r.h.s. are of order **100 GeV**, cancellations are required to make it work
- Measure of **fine-tuning**: sensitivity of m_Z to Lagrangian parameters (a la Barbieri-Giudice, but at the weak scale):

$$\delta(\xi) = \left| \frac{\partial \log m_Z^2}{\partial \log \xi} \right|, \quad \xi = m_u^2, m_d^2, b, \mu \longrightarrow \underline{\mu, \tan \beta, m_A, m_Z}$$

- A useful approximation ($\tan \beta \gg 1$):

$$\delta(\mu) \approx \frac{4\mu^2}{m_Z^2}, \quad \delta(b) \approx \frac{4m_A^2}{m_Z^2 \tan \beta}.$$

Same pars. as the t-channel direct det. diagram!

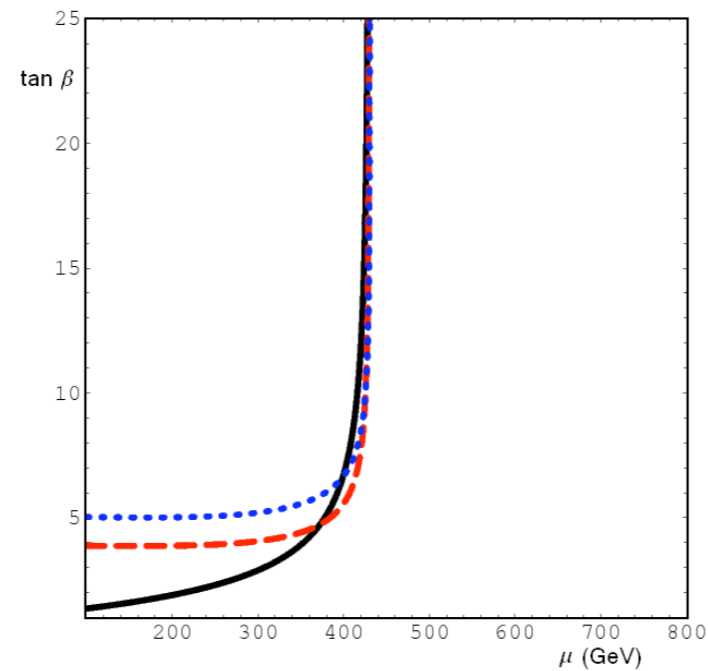


Figure 1: Contours of 1% fine-tuning in the $(\mu, \tan \beta)$ plane. The black (solid) contour corresponds to $m_A = 100$ GeV, but remains essentially unchanged for any value of m_A in the range between 100 and 1000 GeV. The red (dashed) and blue (dotted) contours correspond to $m_A = 1.5$ and 2 TeV, respectively.

[from [MP, Spethmann, hep-ph/0702038](#)]

Fine-Tuning Beyond Tree Level

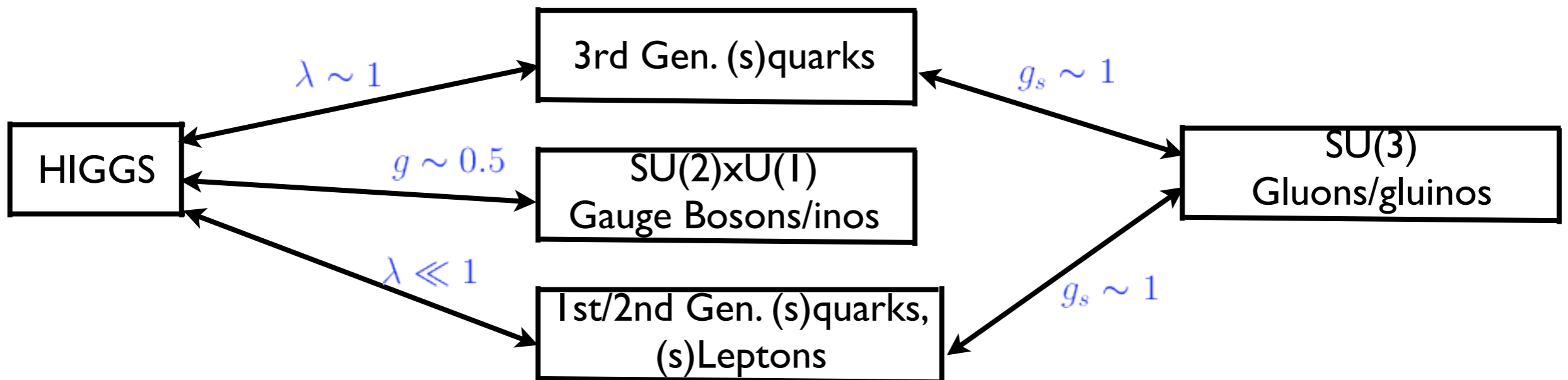
- Consider once again the Z mass formula:

$$m_Z^2 = -m_u^2 \left(1 - \frac{1}{\cos 2\beta}\right) - m_d^2 \left(1 + \frac{1}{\cos 2\beta}\right) - 2|\mu|^2$$

- At loop level, the Higgs mass par's. receive quadratically divergent corrections, cut off by superpartner masses ("SUSY solves the hierarchy problem")

$$\delta m_h^2 \sim \frac{g_p^2}{16\pi^2} \Lambda^2 \quad (\text{SM}) \quad \longrightarrow \quad \delta m_h^2 \sim \frac{g_p^2}{16\pi^2} m_{\tilde{p}}^2 \log \frac{\Lambda^2}{m_{\tilde{p}}^2} \quad (\text{SUSY})$$

- While a large number of parameters enter, the "hierarchy of couplings" in the SM/MSSM simplifies the problem:



Fine-Tuning Beyond Tree Level

- So: 3rd gen. squark loops are the most important quantitatively
- Other superpartners may be a factor of 5 or more heavier than the 3rd gen squarks with no effect on fine-tuning
- Gluino first appears at 2 loops - additional suppression of its effect on fine-tuning
- Assuming moderate $\tan \beta$, top/stop dominates over bottom/sbottom, giving

$$\delta m_{H_u}^2 \approx \frac{3}{16\pi^2} \left(y_t^2 (\tilde{m}_{Q_3}^2 + \tilde{m}_{t^c}^2) + y_t^2 (A_t \sin \beta - \mu \cos \beta)^2 \right) \log \frac{2\Lambda^2}{\tilde{m}_{Q_3}^2 + \tilde{m}_{t^c}^2}$$

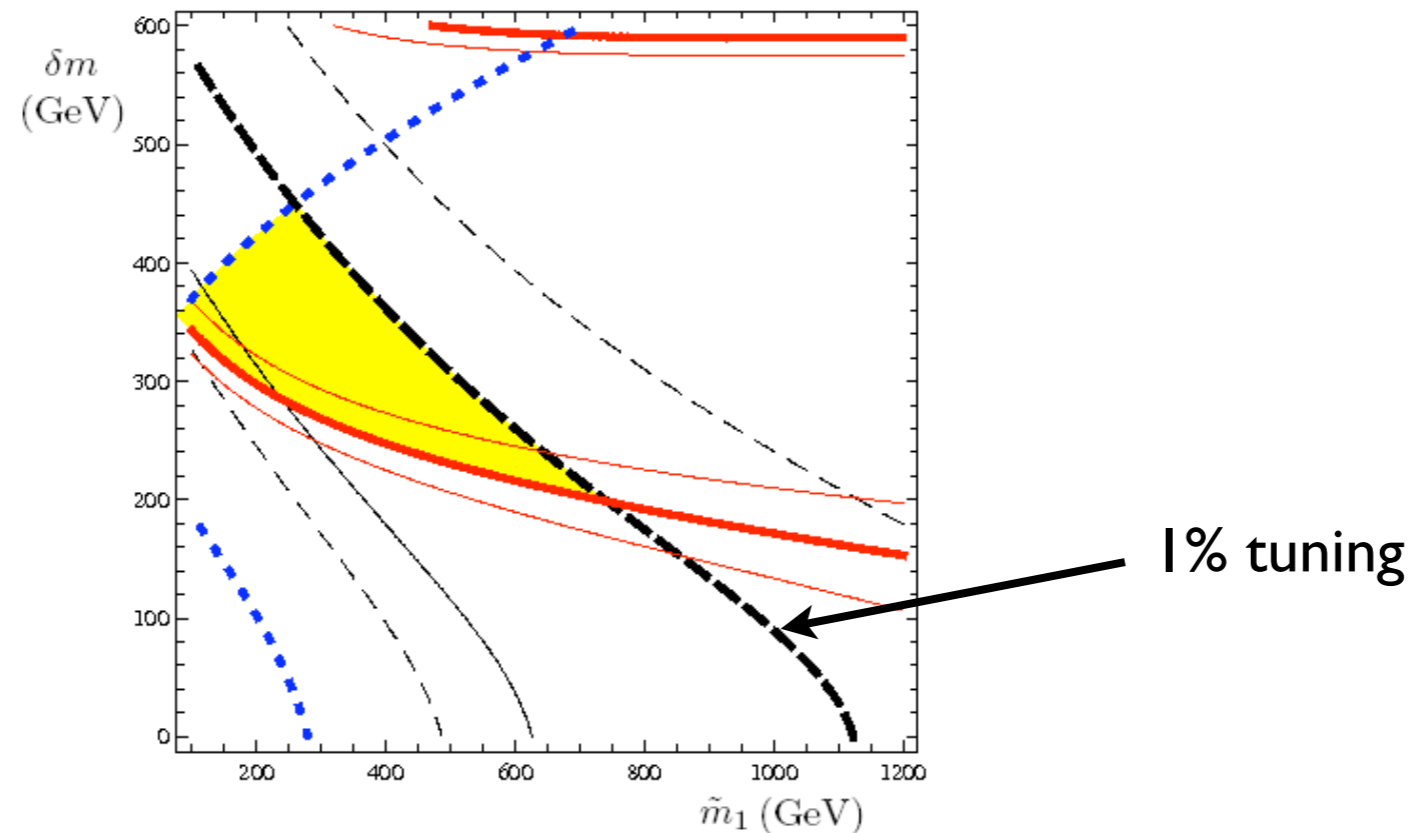
$$\delta_t m_Z^2 \approx -\delta m_{H_u}^2 \left(1 - \frac{1}{\cos 2\beta} \right)$$

- Quantify fine-tuning:

$$\Delta_t = \left| \frac{\delta_t m_Z^2}{m_Z^2} \right|.$$

The “Golden Region”

- At tree level, $m_h > m_Z$, while LEP-2 requires $m_h > 114 \text{ GeV}$
- Solution: large loop corrections \rightarrow tension with fine-tuning! (“little hierarchy problem”)



$$\Lambda = 100 \text{ TeV}, \quad \theta_t = \pi/4, \quad \tan \beta = 10$$

[plot from MP, Spethmann, hep-ph/0702038]

What About the LHC?

BBC NEWS

SCIENCE & ENVIRONMENT

27 August 2011 Last updated at 02:41 ET

LHC results put supersymmetry theory 'on the spot'

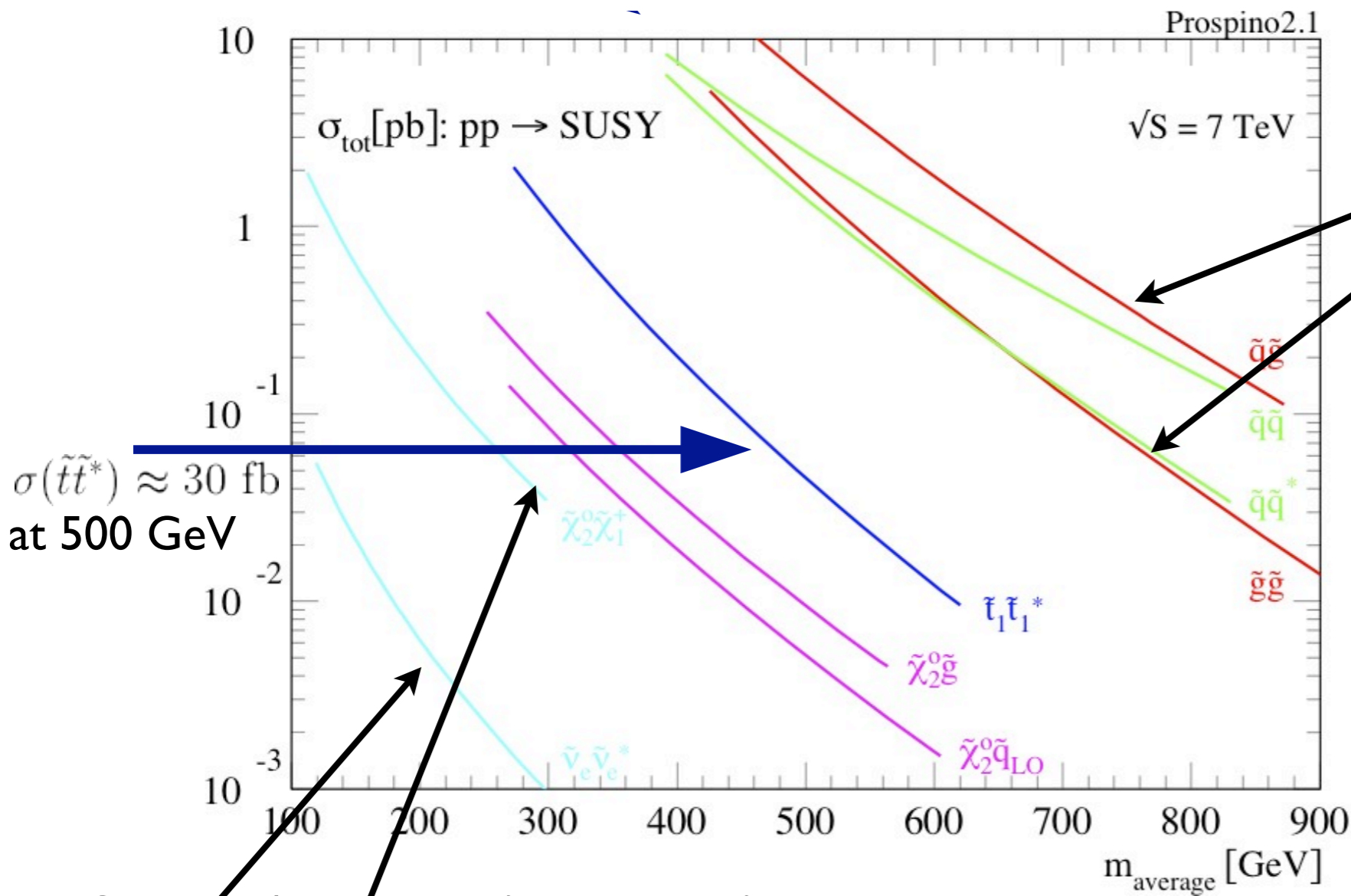


By Pallab Ghosh
Science correspondent, BBC News

Results from the Large Hadron Collider (LHC) have all but killed the simplest version of an enticing theory of sub-atomic physics.

Researchers failed to find evidence of so-called "supersymmetric" particles, which many physicists had hoped would plug holes in the current theory.

What About the LHC?

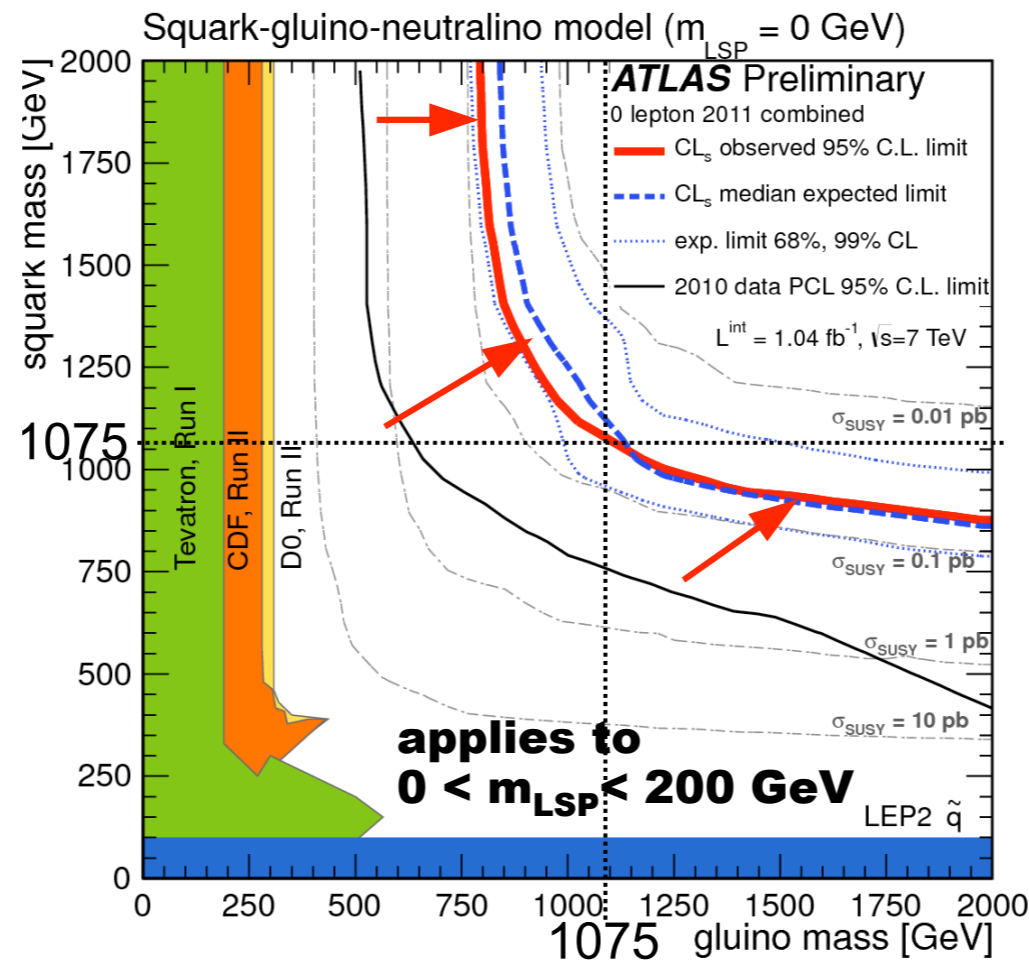


All searches so far rely on producing **gluinos** and/or **1st, 2nd gen. squarks**, different decay channels

Chargino/neutralino (e.g. higgsino) cross sections are also small

Plot credit: H. Bachacou talk at LP-11

What About the LHC?



Plot credit:
 H. Bachacou
 talk at LP-11

BOTTOM LINE: So far the LHC has had **NO*** impact on fine-tuning in the MSSM
 [Not so in specific SUSY breaking models, e.g. where three gen. of squarks have
 common mass term at some scale]

Now, Back to Dark Matter

Procedure: Parameter Scans

- Generate ~ 1000000 random points in the 5-dim. parameter space $(M_1, M_2, \mu, m_A, \tan \beta)$ [all real, $>0!$]
- Uniformly distributed in $\log M_1, \log M_2, \log \mu, \log m_A, \tan \beta$
- Scan boundaries:
$$M_1 \in [10, 10^4] \text{ GeV}; \quad M_2 \in [80, 10^4] \text{ GeV};$$
$$\mu \in [80, 10^4] \text{ GeV}; \quad m_A \in [100, 10^4] \text{ GeV};$$
$$\tan \beta \in [2, 50].$$
- Require: neutral LSP, no charginos below 100 GeV.
Note: we do **not** require correct relic density!
- Fix $m_h = 120 \text{ GeV}$ (assume top/stop loops fix it)

DD Cross Section

$$\sigma = \frac{4m_r^2 f_p^2}{\pi} \quad \frac{f_p}{m_p} = \sum_{q=u,d,s} f_{T_q}^{(p)} A_q + \frac{2}{27} f_{TG}^{(p)} \sum_{q=c,b,t} A_q. \quad f_{TG}^{(p)} = 1 - \sum_{q=u,d,s} f_{T_q}^{(p)}.$$

$$f_{Tu}^{(p)} = 0.08, \quad f_{Td}^{(p)} = 0.037, \quad \underline{f_{Ts}^{(p)} = 0.34}.$$

$$A_i = -\frac{g}{4m_W B_i} \left[\left(\frac{D_i^2}{m_h^2} + \frac{C_i^2}{m_H^2} \right) \text{Re} [\delta_{2i}(gZ_{\chi 2} - g'Z_{\chi 1})] \right. \\ \left. + C_i D_i \left(\frac{1}{m_h^2} - \frac{1}{m_H^2} \right) \text{Re} [\delta_{1i}(gZ_{\chi 2} - g'Z_{\chi 1})] \right], \quad (6)$$

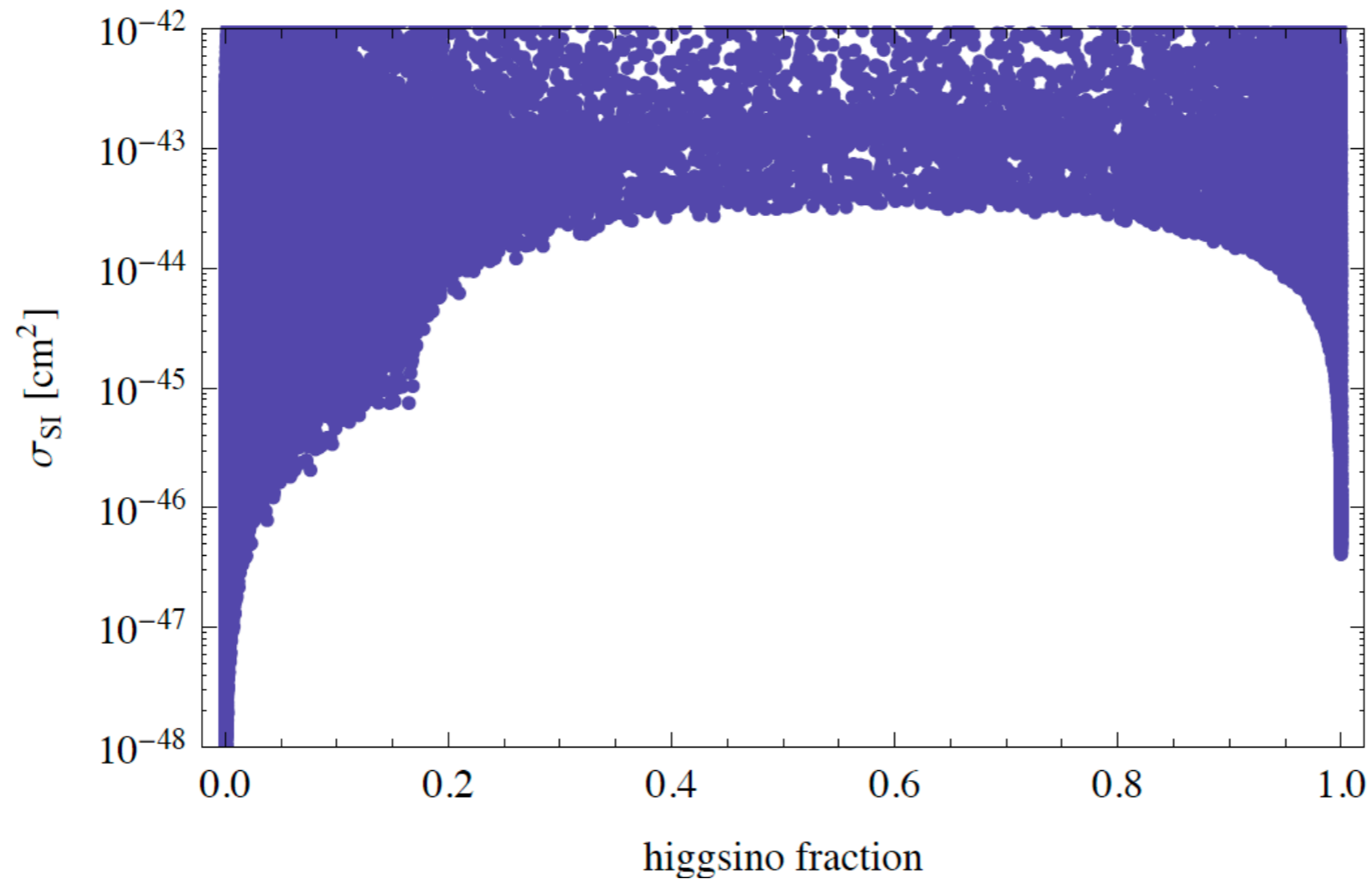
where for up-type quarks

$$B_u = \sin \beta, \quad C_u = \sin \alpha, \quad D_u = \cos \alpha, \quad \delta_{1u} = Z_{\chi 3}, \quad \delta_{2u} = Z_{\chi 4}; \quad (7)$$

while for down-type quarks

$$B_d = \cos \beta, \quad C_d = \cos \alpha, \quad D_d = -\sin \alpha, \quad \delta_{1d} = Z_{\chi 4}, \quad \delta_{2d} = -Z_{\chi 3}. \quad (8)$$

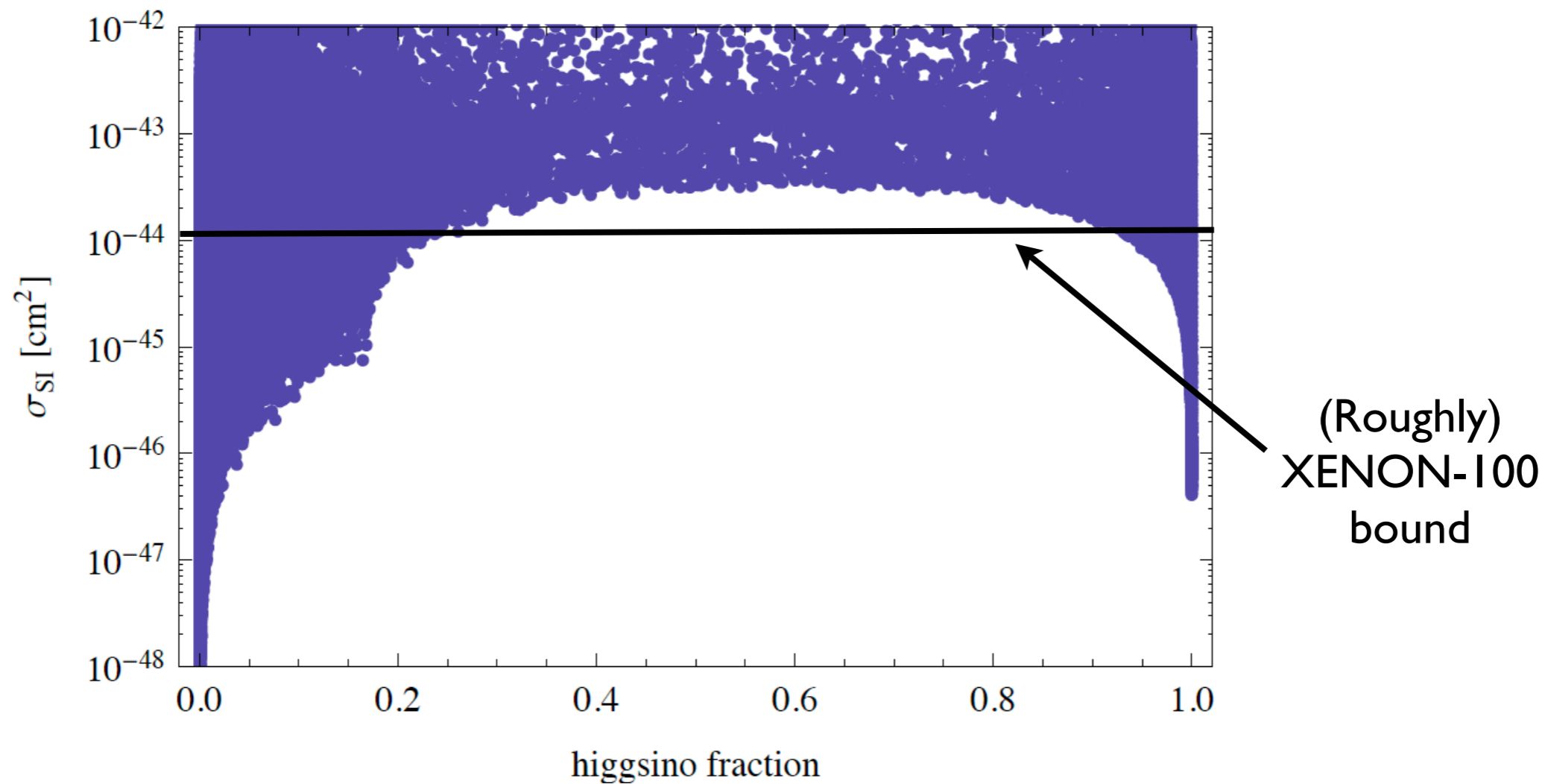
Results: Higgsino Fraction



$$F_H = |Z_{\chi 3}|^2 + |Z_{\chi 4}|^2 .$$

$$\tilde{\chi}_1^0 = Z_{\chi 1} \tilde{B} + Z_{\chi 2} \tilde{W}^3 + Z_{\chi 3} \tilde{H}_d^0 + Z_{\chi 4} \tilde{H}_u^0$$

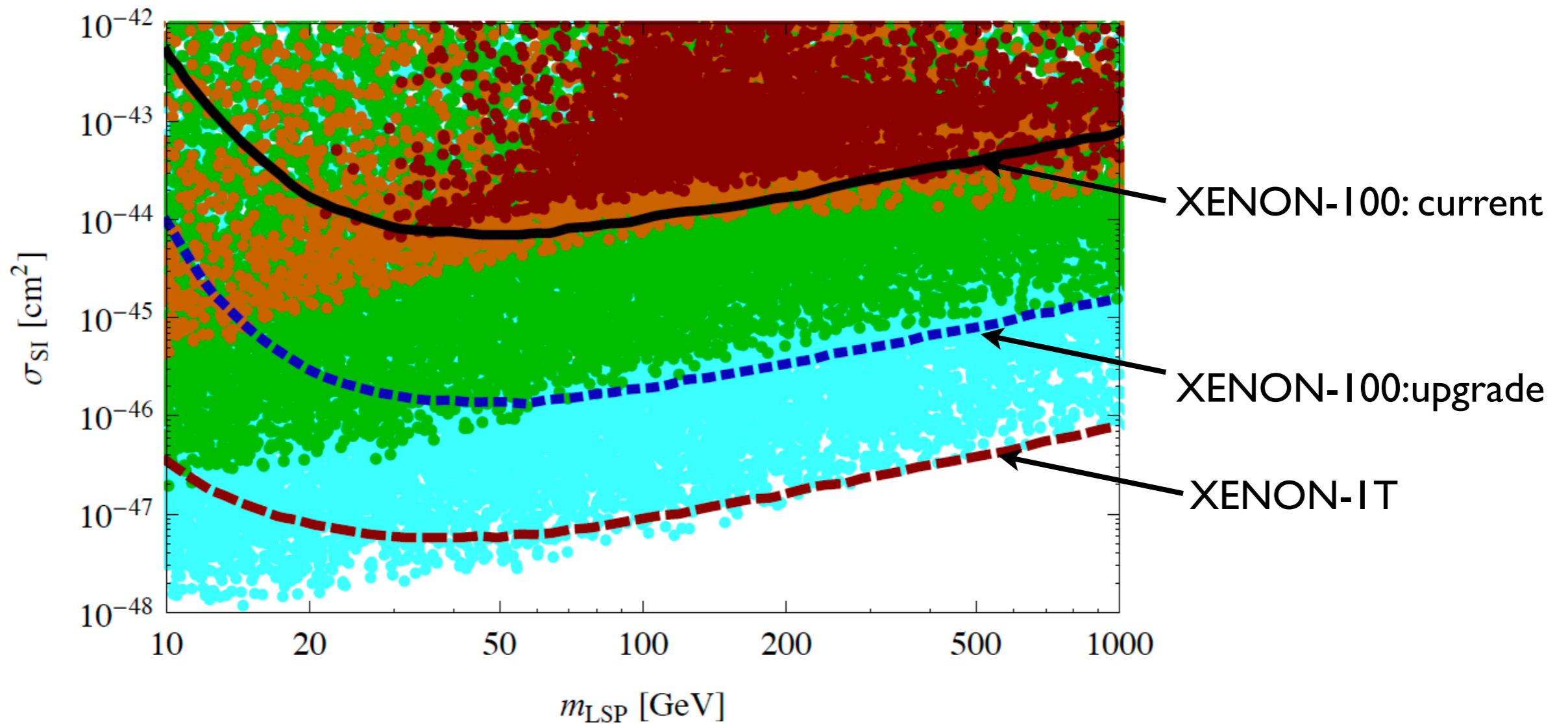
Results: Higgsino Fraction



$$F_H = |Z_{\chi 3}|^2 + |Z_{\chi 4}|^2 .$$

Low DD cross section \longleftrightarrow pure gaugino OR pure Higgsino LSP

Results: Higgsino Fraction

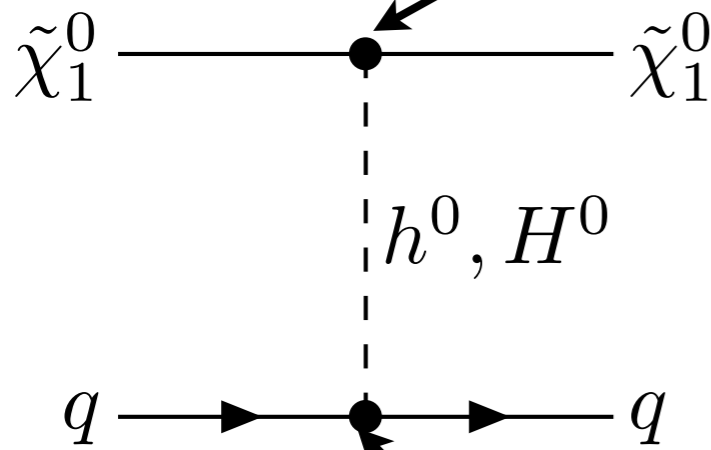


Color-code: “purity” $p = \min(F_H, 1 - F_H)$.

- Red $p > 0.2$
- Orange $0.1 < p < 0.2$
- Green $0.01 < p < 0.1$
- Cyan $0.001 < p < 0.01$

Understanding H-Fraction Bound

$$\begin{aligned} \tilde{\chi}^0 \tilde{\chi}^0 h &: (gZ_{\chi 2} - g'Z_{\chi 1})(\cos \alpha Z_{\chi 4} + \sin \alpha Z_{\chi 3}), \\ \tilde{\chi}^0 \tilde{\chi}^0 H &: (gZ_{\chi 2} - g'Z_{\chi 1})(\sin \alpha Z_{\chi 4} - \cos \alpha Z_{\chi 3}). \end{aligned}$$



“Generic” mixings: $Z_{\chi i} \sim 1$

“Natural” cross section (h exchange only):

$$\sigma \sim (\text{a few}) \times 10^{-44} \text{ cm}^2$$

➡ $Z_{\chi i} \ll 1$ required to get lower x-section!

$$u\bar{u}h : \frac{\sqrt{2}m_u}{v \cos \beta}$$

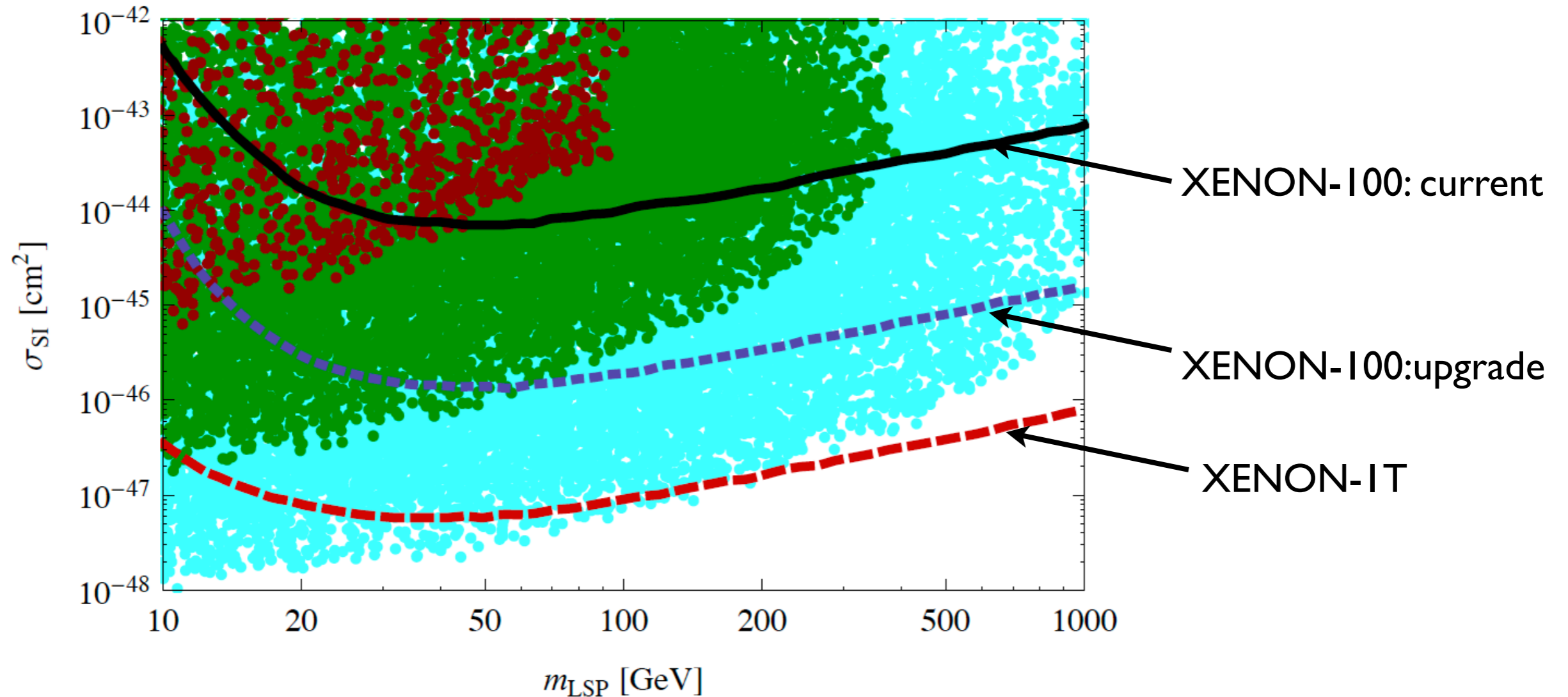
$$d\bar{d}h : \frac{\sqrt{2}m_d}{v \sin \beta}$$

(basically fixed)

Higgsino Fraction Summary

- **Generic** DD cross section from light Higgs exchange is $\sigma \sim (\text{a few}) \times 10^{-44} \text{ cm}^2$, **already being explored** by XENON-100 and others
- **Pure-gaugino** or **pure-higgsino** LSP is required to suppress the cross section below this level
- Current bounds: $p > 0.2$ for any mLSP, and $p > 0.1$ for $m\text{LSP} > 50 \text{ GeV}$, is ruled out
- In the rest of the talk, we consider the “**gaugino LSP sample**”, $M_1 < \mu$ and/or $M_2 < \mu$, and “**Higgsino LSP sample**”, $M_1 > \mu$ and $M_2 > \mu$, separately

EWSB Fine-Tuning: Gaugino LSP



Color-code: EWSB fine-tuning

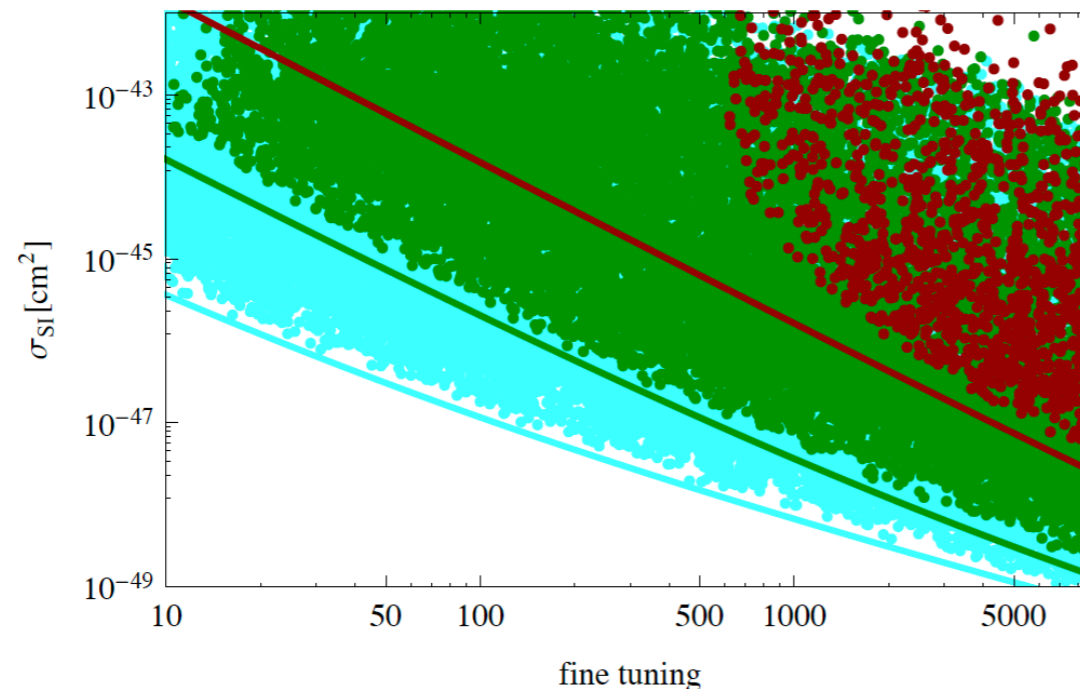
Red $\Delta < 10$
 Green $10 < \Delta < 100$
 Cyan $100 < \Delta < 1000$

Lower DD cross section
means
MORE FINE-TUNING!

EWWSB Fine-Tuning: Gaugino LSP

- This correlation is easy to understand: Smaller DD cross section \Rightarrow purer gaugino LSP \Rightarrow larger μ \Rightarrow more severe EWWSB fine-tuning
- Minimal DD cross section consistent with given fine-tuning*:

$$\sigma_{\min} = (1.2 \times 10^{-42} \text{ cm}^2) \left(\frac{120 \text{ GeV}}{m_h} \right)^4 \frac{1}{\Delta} \left(\frac{1}{\tan \beta} + \frac{1}{\sqrt{\Delta}} \frac{M_{\text{LSP}}}{m_Z} \right)^2$$



Color-code: LSP mass

- Red $M_{\text{LSP}} \in [10, 100] \text{ GeV}$
- Green $M_{\text{LSP}} \in (100, 1000] \text{ GeV}$
- Cyan $M_{\text{LSP}} \in (1, 10] \text{ TeV}$

$$(\tan \beta)_{\max} = 50$$

- * Positive, real parameters are required! (more general case - discuss later)

Gauginino LSP Summary

- Observed an **inverse correlation** between **DD cross section** and **EWSB fine-tuning**
- Obtained a simple **analytic formula** for minimal DD cross section consistent with given FT
- Current bounds already **relevant**:
XENON-100 implies FT **worse than 1/10** for **$m_{LSP} > 70$ GeV**
- XENON-1T will probe down to **FT ~ 1%**!

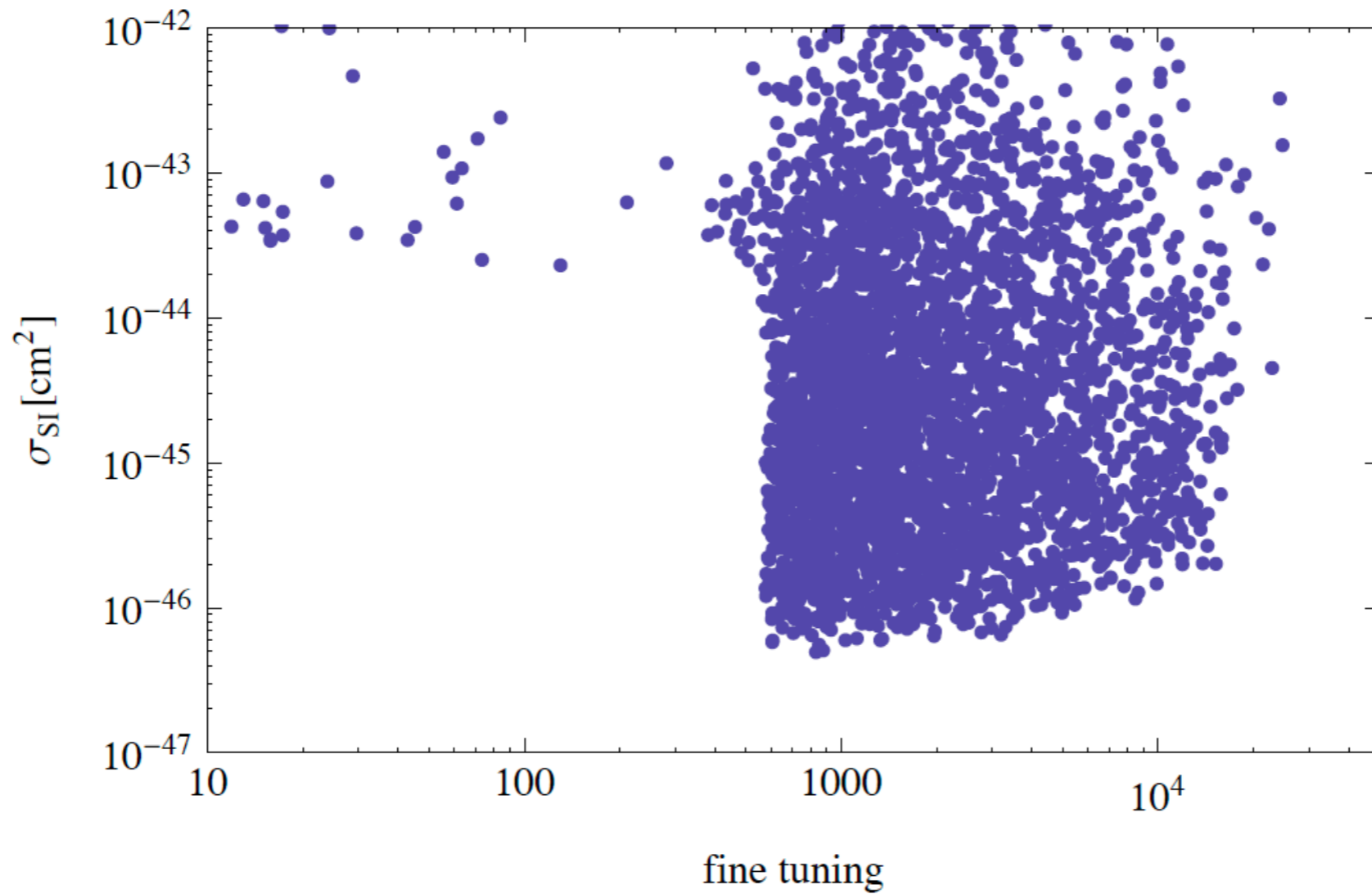
Higgsino LSP

- No correlation between DD cross section and FT in this case: μ fixed, $M_{1,2} \rightarrow \infty \Rightarrow \sigma \rightarrow 0$
- Add a **relic density** constraint: $\Omega_{\text{dm}} h^2 = 0.110 \pm 0.006$
- For fixed MSSM parameters, relic density uniquely **predicted** (assuming standard FRW cosmology up to $T \sim M_{\text{LSP}}$):

$$\Omega h^2 \approx 10^9 \frac{(n+1)x_f^{n+1} \text{ GeV}^{-1}}{(g_{*S}/g_*^{1/2})m_{\text{Pl}}\sigma_0}$$

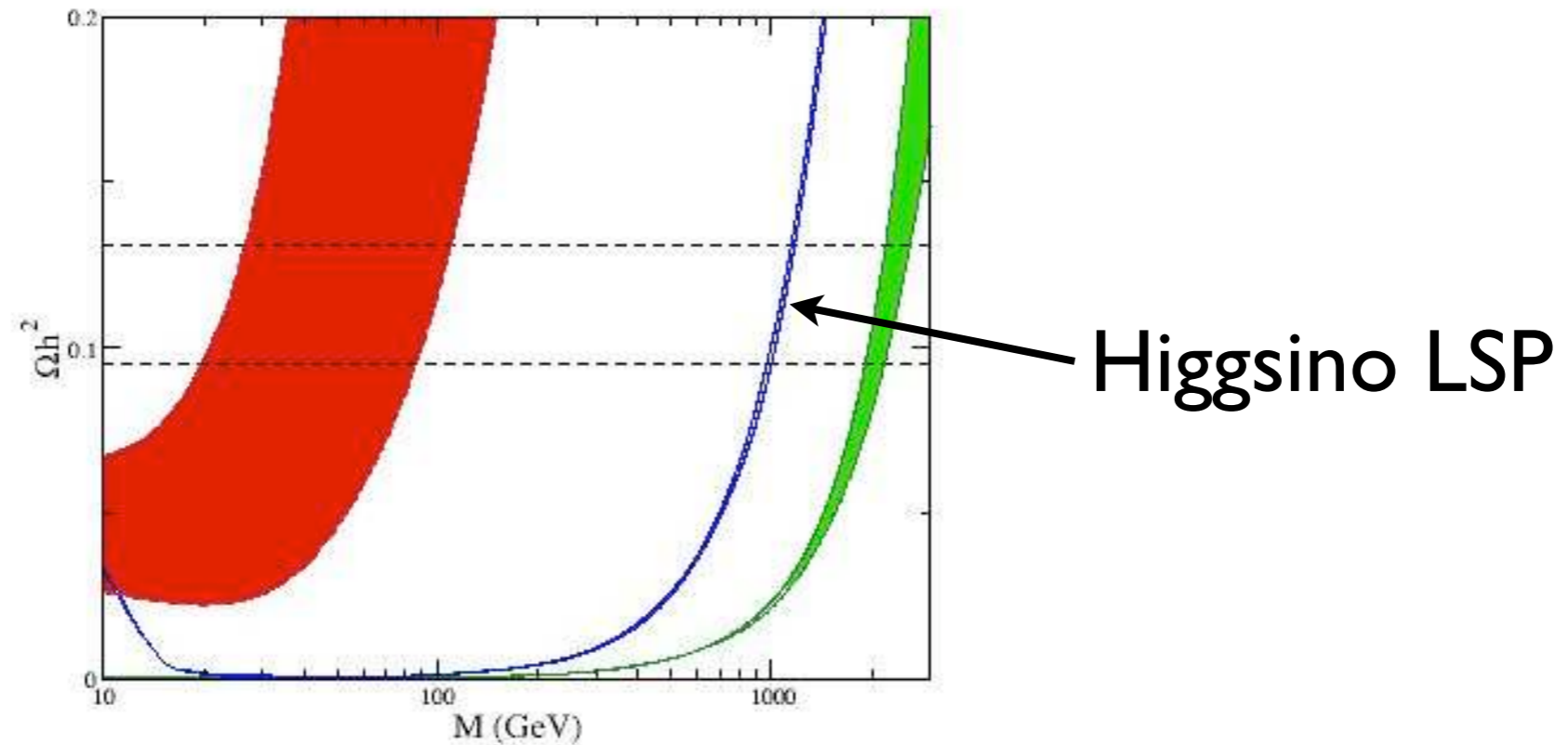
- Higgsino (co-)annihilation channels: **W/Zs, quarks** (predominantly 3rd gen.)
- Our approach: ignore quark final states, impose a **“one-sided”** RD constraint $\Omega_{\text{pred}} \geq \Omega_{\text{obs}}$

Higgsino LSP



Conclusion: XENON-100 already implies $FT > 500$, if the (one-sided) relic density constraint is assumed!

Higgsino LSP Relic Density



Observed (or higher) relic density requires $\mu \geq 1$ TeV \Rightarrow EWSB FINE-TUNING!

[Plot from [Arkani-Hamed, Delgado, Giudice, hep-ph/0601041](#)]

Including Signs/Phases

- While the Higgs/gauge sector of the MSSM has 5 parameters, $(M_1, M_2, \mu, m_A, \tan \beta)$, only **2 phases** are physical: $\varphi_1 = \arg(\mu M_1 \sin 2\beta)$, $\varphi_2 = \arg(\mu M_2 \sin 2\beta)$
- Repeat the scan, allowing **both signs** of M_1, M_2
- New feature: **accidental cancellations** within the t-channel Higgs exchange amplitude are possible, lowering the cross section
- **Quantify** these cancellations: $\Delta_{\text{acc}} \equiv \sqrt{\sum_{i=1}^5 \left(\frac{\partial \log \sigma}{\partial \log p_i} \right)^2}$

Accidental Cancellations

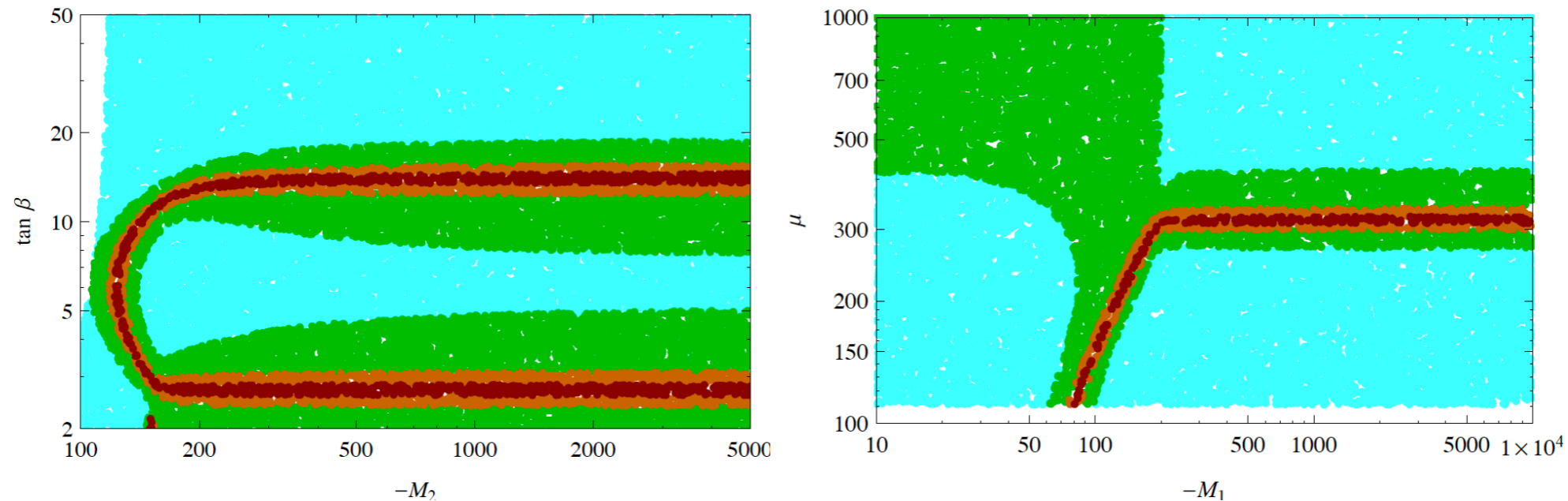
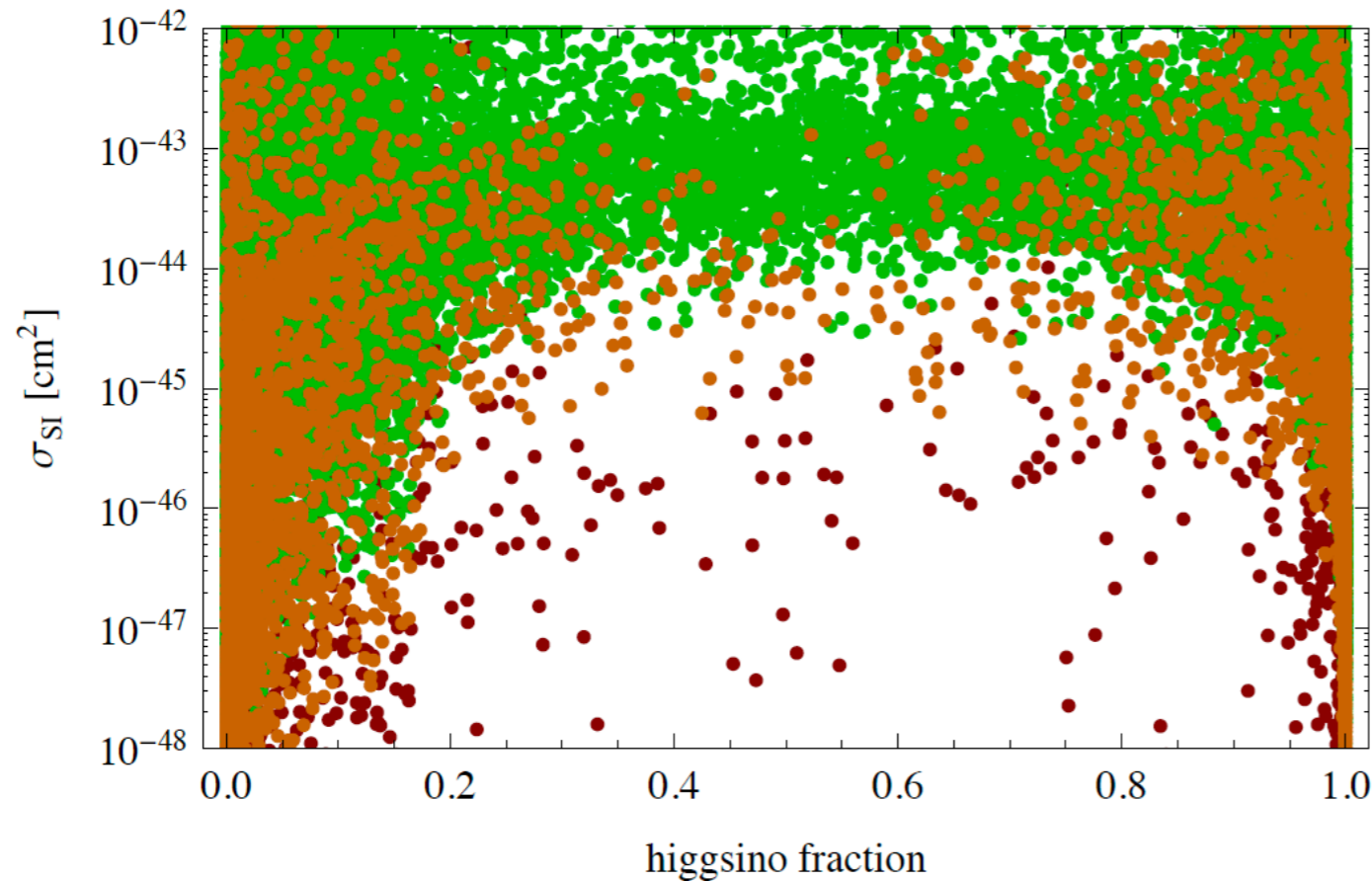


Figure 10: Left panel: Scatter plot of direct detection cross section as a function of $\tan \beta$ and $-M_2$, with $M_1 = -150$ GeV, $\mu = 200$ GeV, and $m_A = 500$ GeV. Cyan, green, orange and red points have $\log_{10} \sigma_{\text{cm}^2} > -45$, $\log_{10} \sigma_{\text{cm}^2} \in (-46, -45)$, $\log_{10} \sigma_{\text{cm}^2} \in (-46, -47)$ and $\log_{10} \sigma_{\text{cm}^2} < -47$, respectively. Right panel: Same, as a function of μ and $-M_1$, with $M_2 = -200$ GeV, $\tan \beta = 10$, and $m_A = 500$ GeV.

Results: Higgsino Fraction

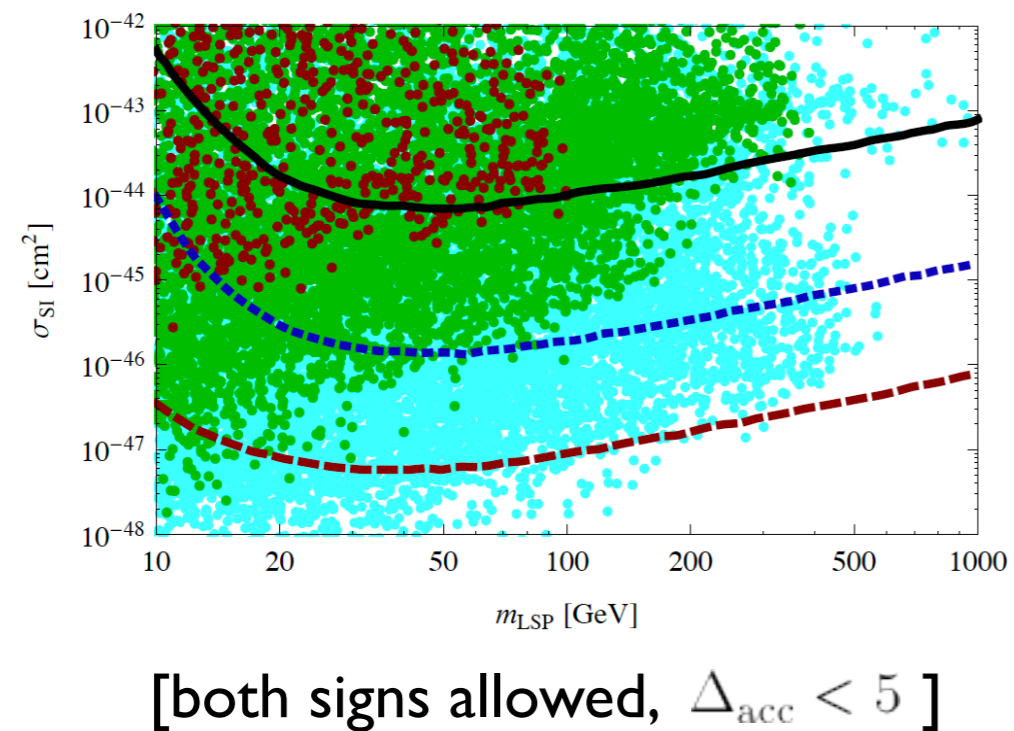
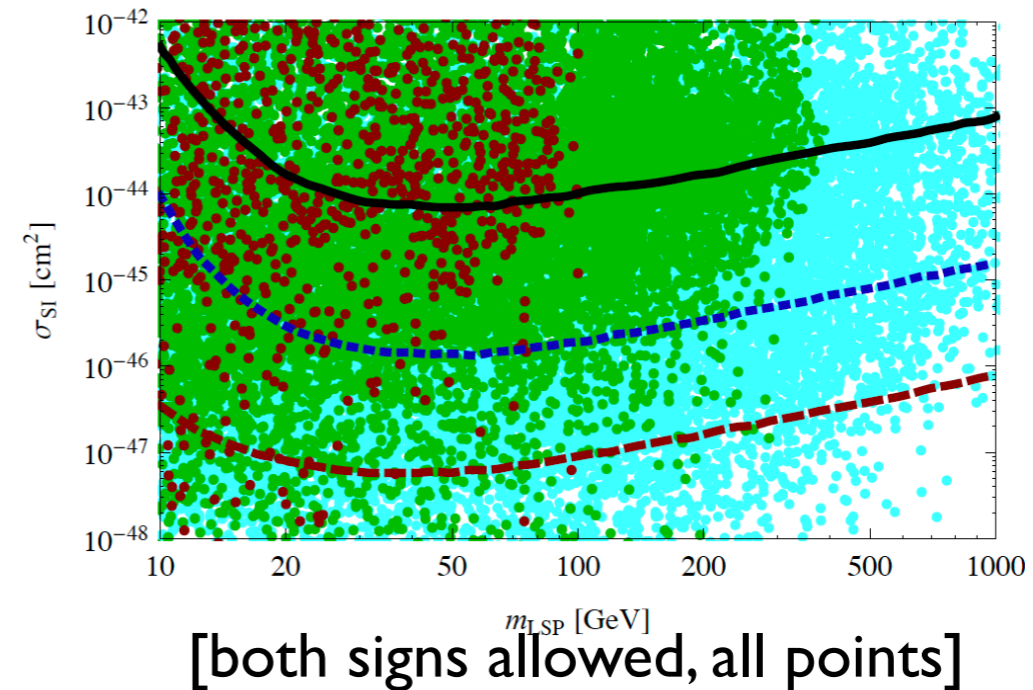
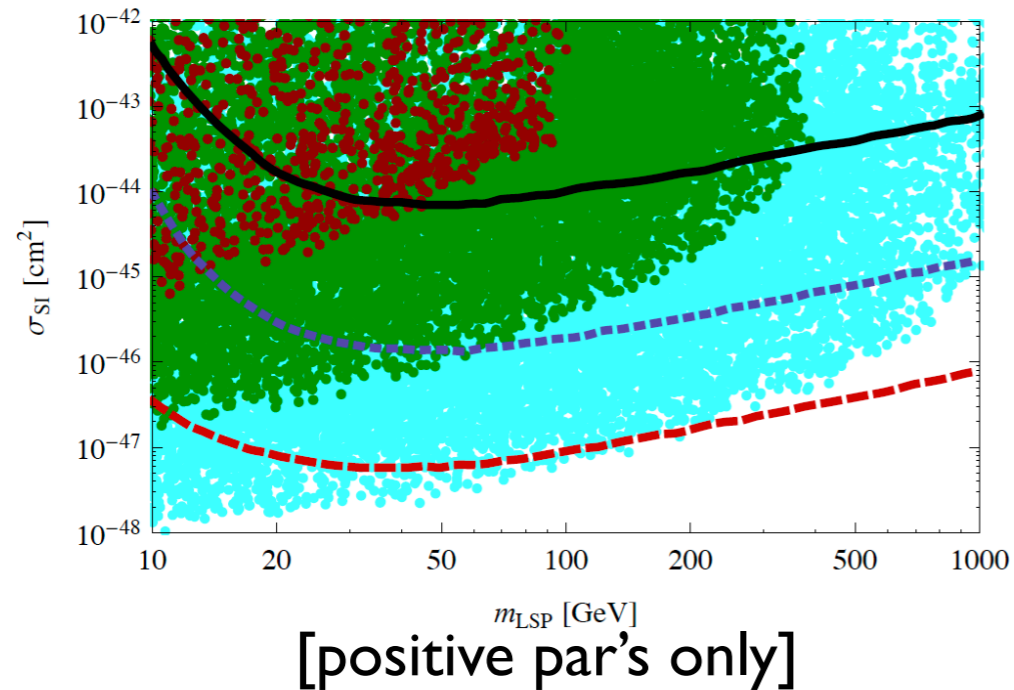


Color-code: "Accidentality"

Red $\Delta_{acc} > 30$
Orange $10 < \Delta_{acc} < 30$
Green $\Delta_{acc} < 10$

Low DD cross section \longleftrightarrow pure gaugino OR pure Higgsino LSP
OR accidental cancellations

EWSB Fine-Tuning for Gaugino LSP



Conclusion: the x-sec/FT correlation still holds, **except** for points with significant accidental cancellations

Conclusions

- “Generic” direct detection cross section (for order-one Higgsino/gaugino mixture LSP) is $\sim 2 \times 10^{-44} \text{ cm}^2$, **ruled out** by XENON-100 for most LSP masses
- Cross sections below 10^{-44} cm^2 require the LSP to be **pure gaugino** or **pure Higgsino**, with lower cross sections corresponding to higher LSP purity
- In the gaugino LSP case, this implies that **lower cross sections** correspond to **stronger fine-tuning** in the EWSB sector
- Current XENON-100 bound implies **FT of 10%** or worse for **$m_{\text{LSP}} > 70 \text{ GeV}$**
- XENON-IT will probe down to **FT~1%** for all LSP masses
- All these statements are true in the **most general MSSM**, assuming only the absence of accidental cancellations in the cross section, do not need to fix thermal relic abundance (e.g. non-standard cosmology is OK)
- In the higgsino LSP case, XENON-100 bound implies **FT of 1/500** or worse **IF** one-sided relic density constraint is imposed